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## MIKHAIL L'VOVICH GOLDMAN

Doctor of physical and mathematical sciences, Professor Mikhail L'vovich Goldman passed away on July 5, 2025, at the age of 80 years.



Mikhail L'vovich was an internationally known expert in science and education. His fundamental scientific articles and text books in various fields of the theory of functions of several variables and functional analysis, the theory of approximation of functions, embedding theorems and harmonic analysis are a significant contribution to the development of mathematics.

Mikhail L'vovich was born on April 13, 1945 in Moscow. In 1963, he graduated from School No. 128 in Moscow with a gold medal and entered the Physics Faculty of the Lomonosov Moscow State University. He graduated in 1969 and became a postgraduate student in the Mathematics Department. In 1972, he defended his PhD thesis "On integral representations and Fourier series of differentiable functions of several variables" under the supervision of Professor Ilyin Vladimir Aleksandrovich, and in 1988, his doctoral thesis "Study of spaces of differentiable functions of several variables with generalized smoothness" at the S.L. Sobolev Institute of Mathematics in Novosibirsk. Scientific degree "Professor of Mathematics" was awarded to him in 1991.

From 1974 to 2000 M.L. Goldman was successively an Assistant Professor, Full Professor, Head of the Mathematical Department at the Moscow Institute of Radio Engineering, Electronics and Automation (technical university). Since 2000 he was a Professor of the Department of Theory of Functions and Differential Equations, then of the S.M. Nikol'skii Mathematical Institute at the Patrice Lumumba Peoples' Friendship University of Russia (RUDN University).

Research interests of M.L. Goldman were: the theory of function spaces, optimal embeddings, integral inequalities, spectral theory of differential operators. Among the most important scientific achievements of M.L. Goldman, we note his research related to the optimal embedding of spaces with generalized smoothness, exact conditions for the convergence of spectral decompositions, descriptions of the integral and differential properties of generalized potentials of the Bessel and Riesz types, exact estimates for operators on cones, descriptions of optimal spaces for cones of functions with monotonicity properties.

M.L. Goldman has published more than 150 scientific articles in central mathematical journals. He is a laureate of the Moscow government competition, a laureate of the RUDN University Prize in Science and Innovation, and a laureate of the RUDN University Prize for supervision of postgraduate students. Under the supervision of Mikhail L'vovich 11 PhD theses were defended. His pupils are actively involved in professional work at leading universities and research institutes in Russia, Kazakhstan, Ethiopia, Rwanda, Colombia, and Mongolia.

Mikhail L'vovich has repeatedly been a guest lecturer and guest professor at universities in Russia, Germany, Sweden, Great Britain, etc., and an invited speaker at many international conferences. Mikhail L'vovich was not only an excellent mathematician and teacher (he always spoke about mathematics and its teaching with great passion), but also a man of the highest culture and erudition, with a deep knowledge of history, literature and art, a very bright, kind and responsive person. This is how he will remain in the hearts of his family, friends, colleagues and students.

The Editorial Board of the Eurasian Mathematical Journal expresses deep condolences to the family, relatives and friends of Mikhail L'vovich Goldman.

# RECONSTRUCTION OF THE WEIGHTED DIFFERENTIAL OPERATOR WITH POINT $\delta$ -INTERACTION

M.Dzh. Manafov, A. Kablan

Communicated by B.E. Kangyzhin

**Key words:** differential operator, uniqueness results, point  $\delta$ -interactions.

**AMS Mathematics Subject Classification:** 34A55, 34B24, 47E05.

**Abstract.** In this article, we provide various uniqueness results for inverse problems of weighted differential operator with point  $\delta$ -interaction. In terms of the method of spectral mappings, we also offer step by step strategies for finding their potential and boundary conditions basing either on the Weyl function, on spectral data, or on two spectra.

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## 1 Introduction

In this work, we study inverse problems for the boundary value problems generated by the differential equation

$$ly := -y'' + q(x)y = \lambda y, \quad x \in (0, a) \cup (a, T) \quad (1.1)$$

with the Robin boundary conditions

$$U(y) := y'(0) - hy(0) = 0, \quad V(y) := y'(T) + Hy(T) = 0, \quad (1.2)$$

and the transmission conditions at the point  $x = a$

$$I(y) := \begin{cases} y(a+0) = y(a-0) \equiv y(a) \\ y'(a+0) - y'(a-0) = -\alpha \lambda y(a), \end{cases} \quad (1.3)$$

where  $q(x)$  is a real function belonging to the space  $L_2[0, T]$ ,  $\lambda$  is a spectral parameter and  $h, H$ , and  $\alpha$  are real numbers with  $\alpha > 0$ . Denote the boundary value problems, defined above, by  $L(q(x), h, H)$ .

It is important to note that, we can interpret problem (1.1) and (1.3) as analyzing the equation

$$-y'' + q(x)y = \lambda \rho(x)y, \quad x \in (0, T), \quad (1.4)$$

when  $\rho(x) = 1 + \alpha \delta(x)$  where  $\delta(x)$  is the Dirac Delta-function (see [1]).

One type of problems, the direct problem, consists of examining the spectral properties of an operator. But some problems in mathematical physics require the investigation of inverse problems of spectral analysis for various differential operators, which require the recovery of operators from some of their given spectral data. Such problems are often considered in mathematics and various branches of natural science and technical science. Direct and inverse problems for the classical Sturm-Liouville operators have been comprehensively investigated in [6, 10, 15] and references therein. Some classes of direct and inverse problems for discontinuous boundary value problems in

various statements have been considered in [2, 7, 8, 12, 13, 16, 17, 18]. Notice that, spectral characteristics for weighted Sturm-Liouville operator with point  $\delta$ -interactions have been investigated in [9, 14]. Here, we provide procedures for finding the potential of a problem and its boundary conditions basing either on the Weyl function, on spectral data, or on two spectra in terms of the method of spectral mappings.

## 2 Constructing the Hilbert space relevant to the problem and some of its spectral properties

We will start this section by defining the Hilbert space  $\mathbb{H} := L_2[0, T] \oplus \mathbb{C}$  of the two component vectors, equipped with the inner product

$$\langle f, g \rangle_{\mathbb{H}} := \int_0^T f_1(x) \bar{g}_1(x) dx + \frac{1}{\alpha} f_2 \bar{g}_2$$

for

$$f = \begin{pmatrix} f_1(x) \\ f_2 \end{pmatrix}, g = \begin{pmatrix} g_1(x) \\ g_2 \end{pmatrix},$$

where  $f_1(x), g_1(x) \in L_2(0, T)$  and  $f_2, g_2 \in \mathbb{C}$ . In the space  $\mathbb{H}$ , we define the operator  $L$

$$L : \mathbb{H} \rightarrow \mathbb{H}$$

with the domain

$$D(L) = \{f \in \mathbb{H} | f_1, f_1' \in AC((0, a) \cup (a, T)), lf_1 \in L_2[(0, T) \setminus \{a\}], f_2 = \alpha f_1(a), U(f_1) = V(f_1) = 0\}$$

and the operator rule

$$L(f) = \begin{pmatrix} lf_1 \\ f_1'(a-0) - f_1'(a+0) \end{pmatrix}.$$

Here,  $AC(\cdot)$  stands for the set of all functions that are absolutely continuous on a related interval.

**Theorem 2.1.** *The operator  $L$  is symmetric.*

*Proof.* We obtain the equality  $\langle Lf, g \rangle_{\mathbb{H}} = \langle f, Lg \rangle_{\mathbb{H}}$  for  $f, g \in D(L)$  immediately from the conditions at the point  $x = a$  and the fact that  $f$  and  $\bar{g}$  satisfy the same boundary conditions (1.2). So,  $L$  is symmetric. □

**Corollary 2.1.** *The function  $W$  defined by  $W\{f, g; x\} = f(x)g'(x) - f'(x)g(x)$  is continuous on  $(0, T)$ .*

**Lemma 2.1.** *If  $y(x, \lambda)$  and  $z(x, \mu)$  are solutions to the equations  $ly = \lambda y$  and  $lz = \mu z$ , respectively, then*

$$\frac{d}{dx} W\{y, z; x\} = (\lambda - \mu)yz.$$

Let  $C(x, \lambda)$ ,  $S(x, \lambda)$ ,  $\varphi(x, \lambda)$  and  $\psi(x, \lambda)$  be solutions to equation (1.1) under the following initial conditions:

$$\begin{aligned} C(0, \lambda) &= 1, \quad C'(0, \lambda) = 0, \quad S(0, \lambda) = 0, \quad S'(0, \lambda) = 1, \\ \varphi(0, \lambda) &= 1, \quad \varphi'(0, \lambda) = h, \quad \psi(T, \lambda) = 1, \quad \psi'(T, \lambda) = -H \end{aligned}$$

and under transmission conditions (1.3). Then,

$$U(\varphi) = V(\psi) = 0.$$

Let us denote

$$\Delta(\lambda) = W\{\varphi, \psi; x\}. \quad (2.1)$$

Due to Corollary 1 and the Ostrogradskii-Liouville theorem (see [4, p. 83])  $W\{\varphi, \psi; x\}$  does not depend on  $x$ . Here, the function  $\Delta(\lambda)$  is called the *characteristic function* of  $L$ . It is easily seen that

$$\Delta(\lambda) = -V(\varphi) = U(\psi), \quad (2.2)$$

and  $\Delta(\lambda)$  is an entire function of  $\lambda$ , so it has at most countable set of zeros  $\{\lambda_n\}_{n \geq 0}$ .

**Lemma 2.2.** *The zeros  $\{\lambda_n\}_{n \geq 0}$  of the characteristic function are the eigenvalues of the boundary value problem  $L$ . Also the functions  $\varphi(x, \lambda_n)$  and  $\psi(x, \lambda_n)$  are the eigenfunctions, and there exists a sequence  $\{\beta_n\}$  such that*

$$\psi(x, \lambda_n) = \beta_n \cdot \varphi(x, \lambda_n), \quad \beta_n \neq 0.$$

Denote

$$\alpha_n := \int_0^T \varphi^2(x, \lambda_n) dx + \alpha \varphi^2(a, \lambda_n). \quad (2.3)$$

The set  $\Omega = \{\lambda_n, \alpha_n\}_{n \geq 0}$  is called the *spectral data* associated with problem (1.1)–(1.3).

**Lemma 2.3.** *The following relation holds*

$$\dot{\Delta}(\lambda_n) = \alpha_n \beta_n,$$

where  $\dot{\Delta}(\lambda) = d\Delta(\lambda)/d\lambda$ .

We omit the proofs of Lemma 2.2 and Lemma 2.3 since they are similar to those for the classical Sturm-Liouville operators (see [11]).

**Corollary 2.2.** *The eigenvalues  $\{\lambda_n\}$  and the eigenfunctions  $\varphi(x, \lambda_n)$ ,  $\psi(x, \lambda_n)$  are real. Also all zeros of  $\Delta(\lambda)$  are simple, i.e.  $\dot{\Delta}(\lambda_n) \neq 0$ .*

Now, consider the solution  $\varphi(x, \lambda)$ . Let  $C_0(x, \lambda)$  and  $S_0(x, \lambda)$  be smooth solutions to equation (1.1) on the interval  $[0, T]$  under the initial condition

$$C_0(x, \lambda) = S'(0, \lambda) = 1, S_0(x, \lambda) = C'_0(0, \lambda) = 0. \quad (2.4)$$

Then,

$$C(x, \lambda) = C_0(x, \lambda), \quad S(x, \lambda) = S_0(x, \lambda), \quad 0 < x < a \quad (2.5)$$

$$\begin{aligned} C(x, \lambda) &= A_1 C_0(x, \lambda) + B_1 S_0(x, \lambda), \\ S(x, \lambda) &= A_2 C_0(x, \lambda) + B_2 S_0(x, \lambda), \quad a < x < T, \end{aligned} \quad (2.6)$$

where

$$\begin{aligned} A_1 &= 1 + \alpha \lambda C_0(a, \lambda) S_0(a, \lambda), \quad B_1 = -\alpha \lambda C_0^2(a, \lambda), \\ A_2 &= \alpha \lambda S_0^2(a, \lambda), \quad B_2 = 1 - \alpha \lambda C_0(a, \lambda) S_0(a, \lambda). \end{aligned} \quad (2.7)$$

Let  $\lambda = \rho^2$ ,  $\rho = \sigma + i\tau$ . It is easy to show that, the function  $C_0(x, \lambda)$  satisfies the following integral equation:

$$C_0(x, \lambda) = \cos \rho x + \frac{1}{\rho} \int_0^x \sin \rho(x-t) q(t) C_0(t, \lambda) dt. \quad (2.8)$$

Using the method of successive approximations to solve problem (2.8), we obtain

$$\begin{aligned} C_0(x, \lambda) &= \cos \rho x + \frac{\sin \rho x}{2\rho} \int_0^x q(t) dt + \frac{1}{2\rho} \int_0^x q(t) \sin \rho(x-2t) dt \\ &\quad + O\left(\frac{1}{\rho^2} \exp(|\tau|x)\right). \end{aligned} \quad (2.9)$$

Analogously,

$$\begin{aligned} S_0(x, \lambda) &= \frac{\sin \rho x}{\rho} - \frac{\cos \rho x}{2\rho^2} \int_0^x q(t) dt + \frac{1}{2\rho^2} \int_0^x q(t) \cos \rho(x-2t) dt \\ &\quad + O\left(\frac{1}{\rho^3} \exp(|\tau|x)\right). \end{aligned} \quad (2.10)$$

By virtue of (2.7) and (2.9)-(2.10),

$$\begin{aligned} A_1 &= \frac{\alpha}{2} \rho \sin 2\rho a - 1 - \frac{\alpha}{2} \cos 2\rho a \int_0^a q(t) dt + O\left(\frac{1}{\rho}\right), \\ B_1 &= -\frac{\alpha}{2} \rho^2 (1 + \cos 2\rho a) - \frac{\alpha}{2} \rho \sin \rho a \int_0^a q(t) dt + O(1), \\ A_2 &= \frac{\alpha}{2} (1 - \cos 2\rho a) + O\left(\frac{1}{\rho}\right), B_2 = -\frac{\alpha}{2} \rho \sin 2\rho a + O(1) \end{aligned}$$

Since  $\varphi(x, \lambda) = C(x, \lambda) + hS(x, \lambda)$ , by using (2.5)-(2.10), we find

$$\varphi(x, \lambda) = \cos \rho x + \left(h + \frac{1}{2} \int_0^x q(t) dt\right) \frac{\sin \rho x}{\rho} + O\left(\frac{1}{\rho} \exp(|\tau|x)\right), 0 < x < a, \quad (2.11)$$

$$\varphi'(x, \lambda) = -\rho \sin \rho x + \left(h + \frac{1}{2} \int_0^x q(t) dt\right) \cos \rho x + O(\exp(|\tau|x)), 0 < x < a, \quad (2.12)$$

$$\begin{aligned} \varphi(x, \lambda) &= \frac{\alpha}{2} \rho (\sin \rho(2a-x) - \sin \rho x) + f_1(x) \cos \rho x + f_2(x) \cos \rho(2a-x) \\ &\quad + O(\exp(|\tau|x)), a < x < T, \end{aligned} \quad (2.13)$$

$$\begin{aligned} \varphi'(x, \lambda) &= -\frac{\alpha}{2} \rho^2 (\cos \rho x + \cos \rho(2a-x)) - \rho f_1(x) \sin \rho x + \rho f_2(x) \sin \rho(2a-x) \\ &\quad + O(\rho \exp(|\tau|x)), a < x < T, \end{aligned} \quad (2.14)$$

where

$$f_1(x) = 1 + \frac{\alpha}{2} h + \frac{\alpha}{4} \int_0^x q(t) dt, \quad f_2(x) = \frac{\alpha}{4} \left(-2h + \int_a^x q(t) dt - \int_0^a q(t) dt\right).$$

It follows from (2.2), (2.13), and (2.14) that

$$\begin{aligned} \Delta(\lambda) &= \frac{\alpha}{2} \rho^2 (\cos \rho T + \cos \rho(2a-T)) + \omega_1 \rho \sin \rho T + \omega_2 \rho \sin \rho(2a-T) \\ &\quad + O(\rho \exp(|\tau|T)), \end{aligned} \quad (2.15)$$

where

$$\begin{aligned} \omega_1 &= -\left(1 + \frac{\alpha}{2} h + \frac{\alpha}{2} H + \frac{\alpha}{4} \int_0^T q(t) dt\right) \\ \omega_2 &= \frac{\alpha}{2} \left(h - H - \frac{1}{2} \int_a^T q(t) dt + \frac{1}{2} \int_0^a q(t) dt\right). \end{aligned}$$

Let  $\lambda_n^0 = (\rho_n^0)^2$  and  $\lambda_n = (\rho_n)^2$  be the zeros of the functions  $\Delta_0(\lambda) = \frac{\alpha}{2}\rho^2(\cos \rho T + \cos \rho(2a - T))$  and  $\Delta(\lambda)$ , respectively.

The following properties of the characteristic function  $\Delta(\lambda)$  and the eigenvalues  $\lambda_n = \rho_n^2$  of the boundary value problem of  $L$  can be discovered using (2.15) and the well-known methods (see, for example [3]).

(i) Denote  $G_\delta = \{\rho : |\rho - \rho_n^0| \geq \delta, n \geq 0\}$ . There is a constant  $C_\delta > 0$  such that

$$|\Delta_0(\lambda)| \geq C_\delta |\lambda| \exp(|\tau|T), \quad \rho \in G_\delta.$$

(ii) For sufficiently large value of  $n$ , the following inequality is valid

$$|\Delta(\lambda) - \Delta_0(\lambda)| \leq \frac{1}{2}C_\delta \exp(|\tau|T), \quad \rho \in \Gamma_n = \{\rho : |\rho| = |\rho_n^0| + \frac{1}{2} \inf_{n \neq m} |\rho_n^0 - \rho_m^0|\}.$$

Thus, for sufficiently large natural number  $n$  and  $\rho \in \Gamma_n$ ,

$$|\Delta_0(\lambda)| \geq C_\delta |\lambda| \exp(|\tau|T) > \frac{1}{2}C_\delta |\lambda| \exp(|\tau|T) > |\Delta(\lambda) - \Delta_0(\lambda)|.$$

Then by Rouché's theorem, the number of zeros of  $\Delta_0(\lambda)$ , counting multiplicities, inside circuit  $G_n$  coincides with the number of zeros of  $\Delta(\lambda)$ . Analogously, applying Rouché's theorem, we say that for sufficiently large values of  $n$ , the function  $\Delta(\lambda)$  has exactly one zero  $\rho_n$  inside each circle  $G_\delta = \{\rho : |\rho - \rho_n^0| \leq \delta\}$ . Since  $\delta$  is arbitrary sufficiently small number, we have

$$\rho_n = \rho_n^0 + \epsilon_n, \quad \epsilon_n = o(1), \quad n \rightarrow \infty \quad (2.16)$$

Since the function  $\Delta_0(\lambda)$  is type of *sinus* (see [5, p. 119]), there exist the number  $d_\delta > 0$  such that, for all  $n$ ,  $|\frac{d}{d\lambda}\Delta_0(\lambda)|_{\lambda=\lambda_n} \geq d_\delta > 0$ . Since  $\rho_n$  are zeros of  $\Delta(\lambda)$ , from (2.15) we get

$$\epsilon_n = -\frac{2}{\alpha \rho_n^0} [\omega_1 \sin \rho_n^0 T + \omega_2 \sin \rho_n^0 (2a - T)] \left[ \frac{d}{d\lambda} \Delta_0(\lambda) \Big|_{\lambda=\lambda_n^0} \right]. \quad (2.17)$$

Substituting (2.17) into (2.16) we get

$$\rho_n = \rho_n^0 + \frac{d_n}{\rho_n^0} + \frac{\gamma_n}{\rho_n^0}, \quad (2.18)$$

where

$$d_n = -\frac{2}{\alpha} [\omega_1 \sin \rho_n^0 T + \omega_2 \sin \rho_n^0 (2a - T)] \left[ \frac{d}{d\lambda} \Delta_0(\lambda) \Big|_{\lambda=\lambda_n^0} \right]$$

and  $\gamma_n = o(1)$ .

Finally, using (2.11), (2.12) and (2.18) into (2.3) we obtain

$$\alpha_n = \alpha_n^0 + o(1), \quad n \rightarrow \infty, \quad (2.19)$$

where

$$\alpha_n^0 = \int_0^T \varphi^2(x, \lambda_n^0) dx + \alpha \varphi^2(a, \lambda_n^0).$$

### 3 Algorithm of solving the inverse problem

In this section, we first give the spectral characteristics of the boundary value problem  $L$  and then demonstrate relationships between their spectral characteristics. Moreover, we provide the formula to solve the inverse problem of the reconstruction of the problem  $L$  basing on the Weyl function, on the spectral data, and on two spectra.

We define the Weyl function by

$$M(\lambda) = \frac{\psi(0, \lambda)}{\Delta(\lambda)}. \quad (3.1)$$

Here the function  $\psi(0, \lambda)$  is the characteristic function of the boundary value problem which consists of equation (1.1) along with the boundary conditions  $y(0) = V(y) = 0$  and transmission conditions (1.3). Let  $\{\mu_n\}_{n \geq 0}$  be the zeros of the entire function  $\psi(0, \lambda)$ . It is clear that,  $\psi(0, \lambda)$  and  $\Delta(\lambda)$  have no common zeros. Thus, the Weyl function  $M(\lambda)$  is meromorphic which has poles at the points  $\{\lambda_n\}_{n \geq 0}$  and zeros at the points  $\{\mu_n\}_{n \geq 0}$ .

The following lemma gives the relationships between the spectral characteristic of  $L$  : the spectral data  $\Omega$ , the Weyl function  $M(\lambda)$  and the two spectra  $\{\lambda_n, \mu_n\}_{n \geq 0}$ .

**Lemma 3.1.** *Let  $M(\lambda)$ ,  $\Omega$  and  $\Delta(\lambda)$  be defined as above. Then the following representation holds:*

$$M(\lambda) = \sum_{h=0}^{\infty} \frac{1}{\alpha_n(\lambda - \lambda_n)}. \quad (3.2)$$

Moreover,  $\Delta(\lambda)$  is uniquely determined up to a multiplicative constant by its zeros:

$$\Delta(\lambda) = T(\lambda_0 - \lambda) \prod_{n=1}^{\infty} \frac{\lambda_n - \lambda}{\lambda_n^0}. \quad (3.3)$$

Since the arguments for proving this lemma are similar to those in [2], we skip the proof. Now, we will consider the following inverse problems of recovering  $L$  :

- *Inverse problem 1:* constructing  $q(x)$ ,  $h$ , and  $H$  when the spectral data  $\{\lambda_n, \alpha_n\}_{n \geq 0}$  is given.
- *Inverse problem 2:* constructing  $q(x)$ ,  $h$  and  $H$  when the Weyl function  $M(\lambda)$  is given.
- *Inverse problem 3:* constructing  $q(x)$ ,  $h$  and  $H$  when the two spectra  $\Omega = \{\lambda_n, \mu_n\}_{n \geq 0}$  are given.

Let us note that, according to (3.1), (3.2) and (3.3), the inverse problems of recovering  $L$  basing on the spectral data and on the two spectra are specifications of the inverse problem of recovering  $L$  from the Weyl function. Consequently, the inverse problems 1–3 are equivalent.

The inverse problems studied here can be seen as generalizations of the inverse problems for the classical Sturm-Liouville operators, see [6, Chapter 1]. In the next section, using results stated above, we provide a constructive procedure for solving these inverse problems.

### 4 Finding solutions to inverse problems

In this section, with the help of Cauchy's integral formula and the Residue theorem, we will solve the inverse problems of recovering the Sturm-Liouville problem  $L(q(x), h, H)$  using the spectrum mappings approach. We first reduce an inverse problem to the so-called main equation which is a

linear equation in a corresponding Banach space of sequences. Finally, we provide the algorithms for solving the inverse problems by using the solution of the main equation.

For this purpose we introduce a new problem with new notations: together with  $L$  we consider a boundary value problem  $\tilde{L}$  of the same form but with different coefficients  $\tilde{q}(x)$ ,  $\tilde{h}$ ,  $\tilde{H}$ . Throughout next sections, if a certain symbol  $e$  denotes an object related to  $L$ , then the symbol  $\tilde{e}$  with tilde denotes the analogous object related to  $\tilde{L}$ . Now we introduce the following notations for convenience of further discussions.

$$\lambda_{n0} = \lambda_n, \quad \lambda_{n1} = \tilde{\lambda}_n, \quad \alpha_{n0} = \alpha_n, \quad \alpha_{n1} = \tilde{\alpha}_n,$$

$$\varphi_{ni}(x) = \varphi(x, \lambda_{ni}), \quad \tilde{\varphi}_{ni}(x) = \tilde{\varphi}(x, \lambda_{ni}),$$

$$Q_{kj}(x, \lambda) = \frac{(\varphi(x, \lambda), \varphi_{kj}(x))}{\alpha_{kj}(\lambda - \lambda_{kj})} = \frac{1}{\alpha_{kj}} \int_0^x \varphi(t, \lambda) \varphi_{kj}(t) dt,$$

$$Q_{ni,kj}(x) = Q_{kj}(x, \lambda_{ni}),$$

for  $i, j = 0, 1$  and  $n, k \geq 0$ . Here  $\tilde{\varphi}(x, \lambda)$  is the solution of (1.4) with the potential  $\tilde{q}$  under the initial conditions  $\tilde{\varphi}(0, \lambda) = 1$ ,  $\tilde{\varphi}'(0, \lambda) = \tilde{h}$ . Similarly, we can define  $\tilde{Q}_{kj}(x, \lambda)$  by replacing  $\varphi$  by  $\tilde{\varphi}$  in the above definition.

Using the Schwartz lemma, see [5, p. 130] and (2.11)-(2.14), (2.17) we obtain the following asymptotic estimates:

$$|\varphi_{ni}(x)| \leq C(|\rho_n^0| + 1), \quad |\varphi_{n0}(x) - \varphi_{n1}(x)| \leq C(|\rho_n^0| + 1)^{\frac{1}{2}}, \quad (4.1)$$

$$|Q_{ni,kj}(x)| \leq \frac{C(|\rho_n^0| + 1)}{(|\rho_n^0 - \rho_k^0| + 1)(|\rho_k^0| + 1)},$$

$$|Q_{ni,k0}(x) - Q_{ni,k1}(x)| \leq \frac{C(|\rho_n| + 1)}{(|\rho_n^0 - \rho_k^0| + 1)(|\rho_k^0| + 1)^{\frac{3}{2}}}, \quad (4.2)$$

$$|Q_{n0,kj}(x) - Q_{n1,j1}(x)| \leq \frac{C(|\rho_n| + 1)^{\frac{1}{2}}}{(|\rho_n^0 - \rho_k^0| + 1)(|\rho_k^0| + 1)}$$

where  $n, k \geq 0, 1$  and  $C > 0$  is independent of  $n, k, i, j$ . Similar estimates are also valid for  $\tilde{\varphi}_{ni}(x)$ ,  $\tilde{Q}_{ni,kj}(x)$ .

**Lemma 4.1.** *Let  $\varphi_{ni}(x)$  and  $Q_{ni,kj}(x)$  be defined as above. Then the following representations hold for  $i, j = 0, 1$  and  $n, k \geq 0$ :*

$$\tilde{\varphi}_{ni}(x) = \varphi_{ni}(x) + \sum_{l=1}^{\infty} (\tilde{Q}_{ni,l0}(x) \varphi_{k0}(x) - \tilde{Q}_{ni,l1}(x) \varphi_{k1}(x)), \quad (4.3)$$

$$Q_{ni,kj}(x) - \tilde{Q}_{ni,kj}(x) + \sum_{l=0}^{\infty} (\tilde{Q}_{ni,l0}(x) Q_{l0,kj}(x) - \tilde{Q}_{ni,l1}(x) Q_{l1,kj}(x)) = 0, \quad (4.4)$$

Both series converge absolutely and uniformly with respect to  $x \in [0, T] \setminus \{a\}$ .



The proof of this lemma is similar to that of the lemma given in [13] and, hence, is omitted.

From all arguments mentioned above, it is seen that, for each fixed  $x \in (0, T) \setminus \{a\}$ , relation (4.3) can be thought as a system of linear equations with respect to  $\varphi_{ni}(x)$  for  $n \geq 0$  and  $i = 0, 1$ . But the series in (4.3) converges only with brackets. So, it is not convenient to use (4.3) as a main equation of the inverse problem. Below, we will transfer (4.3) to a linear equation in a corresponding Banach space of sequences.

Let  $V$  be a set of all indices  $u = (n, i)$ ,  $n \geq 0$ ,  $i = 0, 1$ . For each fixed  $x \in (0, T) \setminus \{a\}$ , we define the vector

$$\phi(x) = [\phi_u(x)] = \begin{bmatrix} \phi_{n0}(x) \\ \phi_{n1}(x) \end{bmatrix}_{n \geq 0}$$

by the formulas

$$\begin{bmatrix} \phi_{n0}(x) \\ \phi_{n1}(x) \end{bmatrix} = \begin{bmatrix} \frac{\rho_n^0+1}{\rho_n} & -\frac{\rho_n^0+1}{\frac{1}{\rho_n^0}} \\ 0 & \frac{1}{\rho_n^0} \end{bmatrix} \begin{bmatrix} \varphi_{n0}(x) \\ \varphi_{n1}(x) \end{bmatrix} = \begin{bmatrix} \frac{(\rho_n^0+1)(\varphi_{n0}(x)-\varphi_{n1}(x))}{\frac{\rho_n}{\varphi_{n1}(x)}} \\ \frac{\varphi_{n1}(x)}{\rho_n} \end{bmatrix} \quad (4.5)$$

Further, we define the block matrix

$$H(x) = [H_{u,v}(x)]_{u,v \in V} = \begin{bmatrix} H_{n0,k0}(x) & H_{n0,k1}(x) \\ H_{n1,k0}(x) & H_{n1,k1}(x) \end{bmatrix}_{n,k \geq 0},$$

where  $u = (n, i)$ ,  $v = (k, j)$  and

$$\begin{bmatrix} H_{n0,k0}(x) & H_{n0,k1}(x) \\ H_{n1,k0}(x) & H_{n1,k1}(x) \end{bmatrix} = \begin{bmatrix} \frac{\rho_n^0+1}{\rho_n} & -\frac{\rho_n^0+1}{\frac{1}{\rho_n^0}} \\ 0 & \frac{1}{\rho_n^0} \end{bmatrix} \begin{bmatrix} Q_{n0,k0}(x) & Q_{n0,k1}(x) \\ Q_{n1,k0}(x) & Q_{n1,k1}(x) \end{bmatrix} \begin{bmatrix} \frac{\rho_k}{\rho_k^0+1} & \rho_k \\ 0 & -\rho_k \end{bmatrix}.$$

Analogously, we define  $\tilde{\varphi}(x)$ ,  $\tilde{H}(x)$  by replacing in the previous definitions  $\varphi_{ni}(x)$ ,  $Q_{ni,kj}(x)$  by  $\tilde{\varphi}_{ni}(x)$ ,  $\tilde{Q}_{ni,kj}(x)$ , respectively. It follows from (2.11) - (2.14), (2.17), (2.18), (4.1), (4.2) and the Schwarz lemma that

$$|\phi_{nj}(x)|, |\tilde{\phi}_{nj}(x)| \leq C, \quad (4.6)$$

and

$$|H_{ni,kj}(x)|, |\tilde{H}_{ni,kj}(x)| \leq \frac{C}{(|\rho_n^0 - \rho_k^0| + 1)(|\rho_k^0| + 1)}, \quad (4.7)$$

where  $C > 0$  is independent of  $n, k, i, j$ .

Let us consider the Banach space  $B$  of bounded sequences  $\alpha = [\alpha_u]_{u \in V}$  with the norm  $\|\alpha\|_B = \sup_{u \in V} |\alpha_u|$  and consider the operator  $E + \tilde{H}(x)$  acting from  $B$  to  $B$ . Here  $E$  is the identity operator. It follows from (4.6), (4.7) that for each fixed  $x$ , this operator is a linear bounded operator, and

$$\|H(x)\|, \|\tilde{H}(x)\| \leq C \sup_n \sum_{k=0}^{\infty} \frac{1}{(|\rho_n^0 - \rho_k^0| + 1)(|\rho_k^0| + 1)} < \infty.$$

Now we are ready to give the main result of this section.

**Theorem 4.1.** *For each fixed  $x \in (0, T) \setminus \{a\}$ , the vector  $\varphi(x) \in B$  satisfies the equation*

$$\tilde{\varphi}(x) = (E + \tilde{H}(x))\varphi(x), \quad (4.8)$$

*in the Banach space  $B$ . Moreover, the operator  $E + \tilde{H}(x)$  has a bounded inverse operator, hence, equation (4.8) is uniquely solvable.*

*Proof.* Using the notation  $\tilde{\phi}(x)$ , rewrite (4.3) as

$$\tilde{\phi}_{ni}(x) = \phi_{ni} + \sum_{n,j} \tilde{H}_{ni,kj}(x) \phi_{kj}(x), (n,i) \in V, (k,j) \in V,$$

which is equivalent to (4.3). Interchanging places for  $L$  and  $\tilde{L}$ , we obtain analogously the equalities

$$\phi(x) = (E - H(x))\tilde{\phi}(x), (E - H(x))(E + \tilde{H}(x)) = E.$$

Hence, the operator  $(E + \tilde{H}(x))^{-1}$  exist, and it is a linear bounded operator.  $\square$

Equation (4.8) is named a *basic equation* of the inverse problem. Solving (4.8) we find the vector  $\phi(x)$ , and hence, the functions  $\varphi_{ni}(x)$ . Thus, we get the following algorithms to find the solution of an inverse problem.

**Algorithm 1.** When the spectral data  $\{\lambda_n, \alpha_n\}_{n \geq 0}$  is given, to construct  $q(x)$ ,  $h$  and  $H$ , we follow the steps:

- (i) first construct  $\tilde{L}$  and then calculate  $\tilde{\phi}(x)$  and  $\tilde{H}(x)$ ,
- (ii) by solving equation (4.8) find  $\phi(x)$  and calculate  $\varphi_{n0}(x)$  by using (4.5),
- (iii) choose some  $n$  (for example,  $n = 0$ ) and construct  $q(x)$ ,  $h$  and  $H$  by using the following formulas:

$$q(x) = \frac{\varphi_{n0}''(x)}{\varphi_{n0}(x)} + \lambda_n, \quad h = \varphi_{n0}'(0), \quad H = -\frac{\varphi_{n0}(T)}{\varphi_{n0}(T)}.$$

**Algorithm 2.** When the function  $M(\lambda)$  is given, to construct  $q(x)$ ,  $h$  and  $H$ , we follow the steps:

- (i) construct the spectral data  $\Omega$  by using (3.2),
- (ii) construct  $q(x)$ ,  $h$  and  $H$  by using Algorithm 1.

**Algorithm 3.** When two spectra  $\{\lambda_n, \mu_n\}_{n \geq 0}$  are given, to construct  $q(x)$ ,  $h$  and  $H$ , we follow the steps:

- (i) calculate  $M(\lambda)$  by using (3.1),
- (ii) construct  $q(x)$ ,  $h$  and  $H$  by using Algorithm 2.

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Manaf Dzh. Manafov  
 Faculty of Arts and Sciences, Department of Mathematics  
 Adıyaman University  
 02040 Adıyaman, Turkey  
 E-mail: mmanafov@adiyaman.edu.tr,

Institute of Mathematics and Mechanics  
 Ministry of Science and Education  
 AZ1141 Baku, Azerbaijan

and

Scientific and Innovation Center  
Western Caspian University  
AZ1001 Baku, Azerbaijan

Abdullah Kablan  
Faculty of Arts and Sciences, Department of Mathematics  
Gaziantep University  
27310 Gaziantep, Turkey  
E-mail: kablan@gantep.edu.tr

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