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#### SHAVKAT ARIFJANOVICH ALIMOV

(to the 75th birthday)



Shavkat Arifjanovich Alimov was born on March 2, 1945 in the city of Nukus, Uzbekistan. In 1968, he graduated from the Department of Mathematics of Physical Faculty of the M.V. Lomonosov Moscow State University (MSU), receiving a diploma with honors. From 1968 to 1970, he was a post-graduate student in the same department under the supervision of Professor V.A. Il'in. He defended his PhD thesis in 1970. In May 1973, at the age of 28, he defended his doctoral thesis devoted to equations of mathematical physics. In 1973, for research on the spectral theory, he was awarded the highest youth prize of the USSR.

From 1974 to 1984, he worked as a professor in the Department of General Mathematics at the Faculty of Computational Mathematics and Cybernetics. In 1984, Sh.A. Alimov joined the Tashkent State University

(TSU) as a professor. From 1985 to 1987 he worked as the Rector of the Samarkand State University, from 1987 to 1990 - the Rector of the TSU, from 1990 to 1992 - the Minister of Higher and Secondary Special Education of the Republic of Uzbekistan. From 1992 to 1994, he headed the Department of Mathematical Physics of the TSU.

After some years of diplomatic work, he continued his academic career as a professor of the Department of Mathematical Physics at the National University of Uzbekistan (NUU). From the first days of the opening of the Tashkent branch of the MSU in 2006, he worked as a professor in the Department of Applied Mathematics. From 2012 to 2017, he headed the Laboratory of Mathematical Modeling of the Malaysian Institute of Microelectronic Systems in Kuala Lumpur. From 2017 to 2019, he worked as a professor at the Department of Differential Equations and Mathematical Physics of the NUU. From 2019 to the present, Sh.A. Alimov is a Scientific Consultant at the Center for Intelligent Software Systems, and an adviser to the Rector of the NUU.

The main scientific activity of Sh.A. Alimov is connected with the spectral theory of partial differential equations and the theory of boundary value problems for equations of mathematical physics. He obtained series of remarkable results in these fields. They cover many important problems of the theory of Schrodinger equations with singular potentials, the theory of boundary control of the heat transfer process, the mathematical problems of peridynamics related to the theory of hypersingular integrals.

In 1984, Sh.A. Alimov was elected a corresponding member and in 2000 an academician of the Academy of Sciences of Uzbekistan. He was awarded several prestigious state prizes.

Sh.A. Alimov has over 150 published scientific and a large number of educational works. Among his pupils there are 10 doctors of sciences and more than 20 candidates of sciences (PhD) working at universities of Uzbekistan, Russia, USA, Finland, and Malaysia.

For about thirty years, Sh.A. Alimov has been actively involved in the reform of mathematical school education.

Sh.A. Alimov meets his 75th birthday in the prime of his life, and the Editorial Board of the Eurasian Mathematical Journal heartily congratulates him on his jubilee and wishes him good health, new successes in scientific and pedagogical activity, family well-being and long years of fruitful life.

#### EURASIAN MATHEMATICAL JOURNAL

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## ON MULTIPERIODIC SOLUTIONS OF PERTURBED NONLINEAR AUTONOMOUS SYSTEMS WITH THE DIFFERENTIATION OPERATOR ON A VECTOR FIELD

#### B.Zh. Omarova, Zh.A. Sartabanov

Communicated by Ya.T. Sultanaev

Key words: multiperiodic solutions, autonomous system, differentiation operator, Lyapunov's vector field, perturbation.

#### AMS Mathematics Subject Classification: 34C46, 35B10, 35C15, 35F35, 35F50.

Abstract. A quasilinear system with the differentiation operator with respect to the directions of vector fields specified by Lyapunov's system with respect to space independent variables and a multiperiodic system with respect to time variables is considered. We study the problem of the existence and uniqueness of a multiperiodic solution of a quasilinear system and we use methods of the theory of multiperiodic solutions of linear systems. The research partially reflects the multiperiodic structure of a solution of the initial problem for quasilinear systems. Conditions for the existence and uniqueness of a multiperiodic solution, an existence theorem of a solution of the initial problem, and the problem of multiperiodic solutions are given. They are proved by the method of contraction mappings defined on spaces of smooth functions.

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## 1 Introduction

Many oscillatory phenomena, such as sound, light, electromagnetic, gas, and hydromechanical phenomena are described by systems of partial differential equations. In this regard, the research on solutions of such systems with oscillatory properties belongs to applied aspects of the theory of ordinary and partial differential equations. Existence conditions of periodic solutions of partial differential equations of the first order with the differentiation operator in the directions of a vector field on a torus were established in [1]. Note that the differential operator of work [1] is similar to the differentiation operator with respect to time variables considered in this paper. We note some similarity of the research methods used in [3] with the methods of this paper. It is explained by the fact that the Poincare-Lyapunov and Hamilton-Jacobi methods are the general basis of the methods for studying such problems.

Comparing the investigations in papers [1, 3] with this work, one should note the common general direction associated with partial differential equations of the first order describing oscillatory processes and some generality of used methods and operators, but the problems investigated are completely different.

The foundations of methods which are used in this paper were laid in [4, 9, 11, 17]. Some results on multiperiodic solutions of systems of equations obtained by the development of the methods of those works are given in  $[2, 5-14]$ , which were further developed in  $[15, 16]$ . In these works, the forces causing systems oscillations were multiperiodic with known periods. Consequently, the authors were looking for solutions with pre-known periods. In the case of autonomous systems, we

do not have such a possibility, but it is possible to reveal multiperiodic structures of solutions of problems accompanying the considered system (e.g. the characteristic equations in time and space variables, corresponding linear systems, etc.), which allows one to determine multiperiodic solutions of the original system.

In this paper, we consider a quasilinear system of equations with the differentiation operator in the directions of vector fields, where the characteristic systems of the differentiation operators with respect to time and space variables are independent. Moreover, the differentiation with respect to space variables is carried out in the directions determined by Lyapunov's system. We studied the problem of the existence of multiperiodic solutions of an analogous autonomous system in our previous works. In this paper, some input data are affected by perturbations depending on time variables. Obviously, multiperiodic solutions of the system are compositions of periodic components with rationally incommensurable periods. Along with multiperiodic components, non-periodic components can also enter a solution of the initial problem. The structure of a solution with selected periodic components is called its multiperiodic structure. The determination of the multiperiodic structure of a solution of the considered system fully reveals the oscillatory nature of the phenomenon described by the system of the initial problem.

In the case of a non-autonomous system, the frequencies of the required multiperiodic oscillations are mainly determined by the system itself. Therefore, the frequencies and their numbers are known in advance.

In the autonomous case, the main difficulty of the problem under consideration is the uncertainty of the frequencies of periodic oscillations, which are components of multiperiodic oscillations determined by a given system. This difficulty was overcome by using the fact that the characteristic vector field satisfies the conditions of Lyapunov's system.

## 2 Integral representation of multiperiodic solutions of a linear system. Statement of the main problem

Let, for  $r > 0$  and  $s \in \mathbb{N}$ ,  $B_r^s$  denote an open ball in  $\mathbb{R}^s$  centered at the origin and let a vector-function  $x = (x_1, \ldots, x_n)$  depending on  $(\tau, t)$ -time and  $\zeta = (\zeta_1, \ldots, \zeta_l)$ -space variables with the components  $\zeta_j = (\xi_j, \eta_j), j = 1, l$  characterize the oscillatory process described by the system of equations

$$
Dx = Ax + f(\tau, t, \zeta, x) \tag{2.1}
$$

with the differentiation operator

$$
D = \frac{\partial}{\partial \tau} + \left\langle a, \frac{\partial}{\partial t} \right\rangle + \left\langle \nu I \zeta + g, \frac{\partial}{\partial \zeta} \right\rangle.
$$
 (2.2)

Here  $\tau \in \mathbb{R}, t = (t_1, \ldots, t_m) \in \mathbb{R}^m, \zeta_j = (\xi_j, \eta_j) \in B_\delta^2, j = \overline{1, l}, \zeta = (\zeta_1, \ldots, \zeta_l) \in B_\delta^{2l}$  $\delta$  > 0;  $\partial$  $\frac{\partial}{\partial t}$  = ∂  $\partial t_1$ , . . . ,  $\partial$  $\partial t_m$  $\setminus$ , ∂  $rac{\delta}{\partial \zeta}$  = ∂  $\partial \zeta_1$ , . . . ,  $\partial$  $\partial \zeta_l$  $\setminus$ , ∂  $\partial \zeta_j$ = ∂  $\partial \xi_j$ ,  $\partial$  $\partial \eta_j$  $\setminus$ are vector differentiation operators;  $\nu = (\nu_1, \ldots, \nu_l)$  is a constant vector,  $\nu I = diag(\nu_1 I_2, \ldots, \nu_l I_2)$ is the matrix with the symplectic unit  $I_2$  of the second order,  $a = (a_1(\tau, t), \ldots, a_m(\tau, t))$  $a(\tau, t)$  and  $g = (g_1(\tau), \ldots, g_l(\tau)) = g(\tau)$  are vector-functions,  $\langle , \rangle$  denotes the scalar product of vectors, A is a constant n-matrix,  $f = f(\tau, t, \zeta, x)$  is an n-vector-function of variables  $(\tau, t, \zeta, x) \in \mathbb{R} \times \mathbb{R}^m \times \overline{B}_{\delta}^{2l} \times \overline{B}_{\Delta}^n.$ 

Assume that the following conditions are satisfied.

a) The vector-function  $a(\tau, t)$  has the property of  $(\theta, \omega)$ -periodicity and smoothness with respect to  $(\tau, t) \in \mathbb{R} \times \mathbb{R}^m$  of order  $(0, e) = (0, 1, \ldots, 1)$ :

$$
a(\tau + \theta, t + q\omega) = a(\tau, t) \in C_{\tau, t}^{(0, e)} (\mathbb{R} \times \mathbb{R}^m), q \in \mathbb{Z}^m,
$$
\n(2.3)

where  $\mathbb{Z}^m$  is the set of all integer vectors  $q = (q_1, \ldots, q_m)$ ,  $q\omega = (q_1\omega_1, \ldots, q_m\omega_m)$ ;  $\omega_0 = \theta, \omega_1, \ldots, \omega_m$  are positive rationally incommensurable periods.

b) The vector coordinates  $\nu_1, \ldots, \nu_l$  are positive rationally incommensurable constants, therefore, numbers  $\alpha_j = 2\pi\nu_j^{-1}$ ,  $j = \overline{1,l}$  are also incommensurable and fulfilled

$$
q_i \alpha_i + q_j \alpha_j \neq 0, q_i^2 + q_j^2 \neq 0, q_i, q_j \in \mathbb{Z}, (i, j = \overline{0, l}).
$$
\n(2.4)

c) The vector-function  $g(\tau) = (g_1(\tau), \ldots, g_l(\tau))$  with the components

$$
g_j(\tau) = (\varphi_j(\tau), \psi_j(\tau)), j = \overline{1, l}
$$

is  $\theta$ -periodic and continuous, i.e.

$$
g(\tau + \theta) = g(\tau) \in C(\mathbb{R}), \qquad (2.5)
$$

moreover, constants  $\theta$  and  $\alpha_j$ ,  $j = 1, l$  also rationally incommensurable:

$$
q_0\theta + q_j\alpha_j \neq 0, q_0^2 + q_j^2 \neq 0, j = \overline{1, l}; q_0, q_j \in \mathbb{Z}.
$$
 (2.6)

Under assumptions  $(2.3)$ – $(2.6)$  on the basis of the vector fields

$$
\frac{dt}{d\tau} = a(\tau, t),\tag{2.7}
$$

$$
\frac{d\zeta}{d\tau} = \nu I \zeta + g(\tau),\tag{2.8}
$$

the characteristics

$$
t = \lambda \left( \tau, \tau^0, t^0 \right), \tag{2.9}
$$

$$
\zeta = Z\left(\tau - \tau^0\right)\left[\zeta^0 - z\left(\tau^0\right)\right] + z\left(\tau\right) \tag{2.10}
$$

are determined, where  $(\tau^0, t^0, \zeta^0) \in \mathbb{R} \times \mathbb{R}^m \times \mathbb{R}^2$ ,  $Z(\tau) = diag[Z_1(\tau), \ldots, Z_l(\tau)], z(\tau) =$  $(z_1(\tau), \ldots, z_l(\tau)), Z_j(\tau) = \begin{pmatrix} \cos \nu_j \tau & -\sin \nu_j \tau \\ \sin \nu_j \tau & \cos \nu_j \tau \end{pmatrix}$  $\sin \nu_j \tau$   $\cos \nu_j \tau$  $\setminus$  $z_j(\tau + \theta) = z_j(\tau),$ 

$$
z_j(\tau) = \left[ Z_j^{-1}(\tau + \theta) - Z_j^{-1}(\tau) \right]^{-1} \int_{\tau}^{\tau + \theta} Z_j^{-1}(s) g_j(s) \, ds, j = \overline{1, l}.
$$

Based on relations (2.9) and (2.10), we define the first integrals of vector fields (2.7) and (2.8) in the form

$$
t^{0} = \lambda \left(\tau^{0}, \tau, t\right), \tag{2.11}
$$

$$
\zeta^{0} = Z\left(\tau^{0} - \tau\right)\left[\zeta - z\left(\tau\right)\right] + z\left(\tau^{0}\right) \equiv \mu\left(\tau^{0}, \tau, \zeta\right),\tag{2.12}
$$

where the following properties hold:

$$
D\lambda(\tau^0, \tau, t) = 0, \lambda(\tau^0, \tau^0, t) = t,
$$
  
\n
$$
\lambda(\tau', \tau'', \lambda(\tau'', \tau, t)) = \lambda(\tau', \tau, t),
$$
  
\n
$$
\lambda(\tau^0 + \theta, \tau + \theta, t + q\omega) = \lambda(\tau^0, \tau, t) + q\omega;
$$
\n(2.13)

$$
D\mu(\tau^0, \tau, \zeta) = 0, \mu(\tau^0, \tau^0, \zeta) = \zeta,
$$
  
\n
$$
\mu(\tau', \tau'', \mu(\tau'', \tau, \zeta)) = \mu(\tau', \tau, \zeta),
$$
  
\n
$$
\mu(\tau^0 + \theta, \tau + \theta, \zeta) = \mu(\tau^0, \tau, \zeta).
$$
\n(2.14)

Also, note that the matrices  $Z_j(\tau + \alpha_j) = Z_j(\tau), j = \overline{1,l}$  are periodic with the incommensurable periods  $\alpha_j$ ,  $j = 1, l$ . Consequently, the matrix  $Z(\tau)$  is quasiperiodic with the frequency basis  $\nu =$  $(\nu_1, \ldots, \nu_l)$ , which is represented by the multiperiodic blocks  $Z_i(\tau)$  with the period vector  $\alpha =$  $(\alpha_1, \ldots, \alpha_l)$ :

$$
Z(\hat{\tau} + q\alpha) = diag[Z_1(\tau_1 + q_1\alpha_1), \dots, Z_l(\tau_l + q_l\alpha_l)]
$$
  
= diag[Z\_1(\tau\_1), \dots, Z\_l(\tau\_l)] = Z(\hat{\tau}), q \in \mathbb{R}^l, \hat{\tau} = (\tau\_1, \dots, \tau\_l) \in \mathbb{R} \times \dots \times \mathbb{R} = \mathbb{R}^l.

In conclusion, we note that by virtue of relations  $(2.11)$ – $(2.14)$  with an arbitrary smooth  $\text{vector-function} \quad u^0(t,\zeta) \quad \in \quad C_{t,\zeta}^{(e,\tilde{e})}(\mathbb{R}^m \times \mathbb{R}^{2l}), \quad \text{we can determine the zero} \quad u(\tau^0,\tau,t,\zeta) = 0.$  $u^0(\lambda(\tau^0,\tau,t),\mu(\tau^0,\tau,\zeta))$  of the operator D:  $Du(\tau^0,\tau,t,\zeta)=0$ , where e and  $\tilde{e}$  are m and 2l-vectors with the unit coordinates.

It is easy to verify that the initial problem for the linear system

$$
Dx = Ax \tag{2.15}
$$

with the initial condition

$$
x|_{\tau=\tau^0} = u(t,\zeta) \in C_{t,\zeta}^{(e,\tilde{e})} \left( \mathbb{R}^m \times \mathbb{R}^{2l} \right)
$$
 (2.16)

is determined by the formula

$$
x\left(\tau^{0},\tau,t,\zeta\right)=X\left(\tau-\tau^{0}\right)u\left(\lambda\left(\tau^{0},\tau,t\right),\mu\left(\tau^{0},\tau,\zeta\right)\right),\tag{2.17}
$$

where  $X(\tau)$  is the matricant of system (2.15):  $DX(\tau) = AX(\tau)$ ,  $X(0) = E$  is the identity matrix.

We note that the condition

$$
det [X (\theta) - E] \neq 0 \tag{2.18}
$$

is necessary for system (2.15) to have only zero  $(\theta, \omega)$ -periodic with respect to  $(\tau, t)$  solution.

It can be proved that if condition  $(2.18)$  is satisfied, then problem  $(2.15)-(2.16)$  has a nonzero  $(\theta, \omega)$ -periodic solution (2.17) if and only if the fundamental system

$$
u\left(\lambda\left(\tau^0, \tau^0 + \theta, t, \mu\left(\tau^0, \tau^0 + \theta, \zeta\right)\right)\right)
$$
  
= 
$$
\left[E - X(\theta)\right]^{-1}\left[u\left(\lambda\left(\tau^0, \tau^0 + \theta, t\right), \mu\left(\tau^0, \tau^0 + \theta, \zeta\right)\right) - u\left(t, \zeta\right)\right]
$$
 (2.19)

has a  $\omega$ -periodic with respect to t nontrivial smooth solution  $u(t,\zeta)$ .

Therefore, if, along with condition (2.18), system (2.19) has only zero solution in the class of functions

$$
u(t + q\omega, \zeta) = u(t, \zeta) \in C_{t, \zeta}^{(e, \tilde{e})} (\mathbb{R}^m \times \mathbb{R}^{2l}), \qquad (2.20)
$$

then system (2.15) does not have  $(\theta, \omega)$ -periodic solutions, except for the zero one.

Now we consider the inhomogeneous system

$$
Dx = Ax + f(\tau, t, \zeta) \tag{2.21}
$$

with the vector-function

$$
f\left(\tau+\theta,t+q\omega,\zeta\right) = f\left(\tau,t,\zeta\right) \in C_{\tau,t,\zeta}^{(0,e,\tilde{e})}\left(\mathbb{R}\times\mathbb{R}^m\times\mathbb{R}^{2l}\right), q\in\mathbb{Z}^m. \tag{2.22}
$$

Suppose that under condition (2.18), system (2.19) has only zero solution from class (2.20). In other words, system (2.15) has only zero  $(\theta, \omega)$ -periodic with respect to  $(\tau, t)$  solution.

Further, in order to write out an integral representation of the  $(\theta, \omega)$ -periodic solution of system (2.21), we use vector-function (2.22) along with the characteristics  $(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta)),$  $(s, \lambda(s, \tau + \theta, t), \mu(s, \tau + \theta, \zeta))$  and the following vector-function

$$
f_{\theta}\left(s,\lambda(s,\tau,t),\mu(s,\tau,\zeta)\right) = \begin{cases} f\left(s,\lambda(s,\tau,t),\mu(s,\tau,\zeta)\right), \tau \le s \le 0, \\ f\left(s,\lambda(s,\tau+\theta,t),\mu(s,\tau+\theta,\zeta)\right), \\ 0 < s \le \tau + \theta. \end{cases} \tag{2.23}
$$

Then under the above conditions, the  $(\theta, \omega)$ -periodic solution  $x^*(\tau, t, \zeta)$  of system (2.21) has the following integral representation

$$
x^*(\tau, t, \zeta) = \left[X^{-1}(\tau + \theta) - X^{-1}(\tau)\right]^{-1} \int_{\tau}^{\tau + \theta} X^{-1}(s) f_{\theta}(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta)) ds, \tag{2.24}
$$

where  $f_{\theta}(s, \lambda(s,\tau,t), \mu(s,\tau,\zeta))$  has form  $(2.23)$ .

Note that vector-function (2.24) is a  $(\theta, \omega)$ -periodic solution of system (2.21) which can be verified by direct calculations. The uniqueness follows from the fact that homogeneous system (2.15) has only zero  $(\theta, \omega)$ -periodic solution.

Now we state the main goal of this paper: investigation of the problem of the existence of a unique  $(\theta, \omega)$ -periodic with respect to  $(\tau, t)$  solution of quasilinear system (2.1) based on the above information on multiperiodic solutions of linear system (2.21).

# 3 Multiperiodic solution of a quasilinear system with the differentiation operator D

We consider quasilinear system  $(2.1)$  under the following assumptions.

- 1 0 . All assumptions with respect to the input data of linear part of system (2.1) remain valid.
- 2<sup>0</sup>. A vector-function  $f(\tau, t, \zeta, x)$  satisfies the condition

$$
f(\tau + \theta, t + q\omega, \zeta, x) = f(\tau, t, \zeta, x) \in C_{\tau, t, \zeta, x}^{(0, e, \tilde{e}, \hat{e})} \left( \mathbb{R} \times \mathbb{R}^m \times \overline{B}_{\delta}^{2l} \times \overline{B}_{\Delta}^n \right),
$$
(3.1)

where  $e, \tilde{e}, \tilde{e}$  are the vectors with the unit components of dimensions  $m, 2l, n$ , respectively.

3<sup>0</sup>. Linear system (2.15) has no  $(\theta, \omega)$ -periodic solutions except zero.

Let  $\theta, \omega, \delta, \gamma > 0$ . We introduce the space  $S^{\theta, \omega}_{\delta, \gamma}$  of all *n*-vector-functions  $x(\tau, t, \zeta)$  which are  $(\theta,\omega)$ -periodic with respect to  $(\tau,t) \in \mathbb{R} \times \mathbb{R}^m$ , continuously differentiable with respect to  $(\tau,t,\zeta) \in$  $\mathbb{R}\times\mathbb{R}^m\times B_\delta^{2l}$ , continuous with the partial derivatives of order one on the closure  $\mathbb{R}\times\mathbb{R}^m\times\overline{B}_\delta^{2l}$  $\frac{a}{\delta}$  and such that  $||x|| < \gamma$ , where

$$
||x|| = ||x||_{\circ} + \sum_{j=0}^{m} \left\| \frac{\partial x}{\partial t_j} \right\|_{\circ} + \sum_{k=1}^{l} \left( \left\| \frac{\partial x}{\partial \xi_k} \right\|_{\circ} + \left\| \frac{\partial x}{\partial \eta_k} \right\|_{\circ} \right)
$$

is the norm of *n*-vector-functions  $x(\tau, t, \zeta)$  of variables  $\tau = t_0, t = (t_1, \ldots, t_m)$ ,  $\zeta = (\zeta_1, \ldots, \zeta_l) = ((\xi_1, \eta_1), \ldots, (\xi_l, \eta_l)),$  where  $||x||_{\circ} = \sup |x(\tau, t, \zeta)|$ , the supremum is taken with respect to  $(\tau, t, \zeta) \in \mathbb{R} \times \mathbb{R}^m \times \overline{B}_{\delta}^{2l}$  $\delta^{2l}$  and  $|\cdot|$  is the Euclidean norm in  $\mathbb{R}^n$ .

We consider the integral operators  $T, P'_k, P''_k, Q_j, R$  with the kernel  $K = K(\tau, s)$  of the form

$$
(Tx)(\tau, t, \zeta) = \int_{\tau}^{\tau+\theta} K(\tau, s) f_{\theta}(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta), x(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta))) ds,
$$
(3.2)

$$
K(\tau, s) = \left[X^{-1}(\tau + \theta) - X^{-1}(\tau)\right]^{-1} X^{-1}(s),\tag{3.3}
$$

$$
(P'_k x)(\tau, t, \zeta) = \int\limits_{\tau}^{\tau + \theta} K(\tau, s) \left[ \frac{\partial f_{\theta}}{\partial \zeta} + \frac{\partial f_{\theta}}{\partial x} \cdot \frac{\partial x}{\partial \zeta} \right] \frac{\partial \zeta}{\partial \zeta_k} \cdot \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \xi_k} ds, \tag{3.4}
$$

$$
(P_k''x)(\tau, t, \zeta) = \int\limits_{\tau}^{\tau+\theta} K(\tau, s) \left[ \frac{\partial f_\theta}{\partial \zeta} + \frac{\partial f_\theta}{\partial x} \cdot \frac{\partial x}{\partial \zeta} \right] \frac{\partial \zeta}{\partial \zeta_k} \cdot \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \eta_k} ds, \tag{3.5}
$$

$$
(Q_j x)(\tau, t, \zeta) = \int\limits_{\tau}^{\tau + \theta} K(\tau, s) \left[ \frac{\partial f_{\theta}}{\partial t_j} + \frac{\partial f_{\theta}}{\partial x} \cdot \frac{\partial x}{\partial t} \right] \frac{\partial \lambda(s, \tau, t)}{\partial t_j} ds, \tag{3.6}
$$

$$
(Rx)(\tau, t, \zeta) = f(\tau, t, \zeta, x) + \int_{\tau}^{\tau + \theta} K(\tau, s) \left[ \frac{\partial f_{\theta}}{\partial \zeta} + \frac{\partial f_{\theta}}{\partial x} \cdot \frac{\partial x}{\partial \zeta} \right] \frac{\partial \mu(s, \tau, \zeta)}{\partial \tau} ds
$$
  
+ 
$$
A(Tx)(\tau, t, \zeta) + \int_{\tau}^{\tau + \theta} K(\tau, s) \left[ \frac{\partial f_{\theta}}{\partial t} + \frac{\partial f_{\theta}}{\partial x} \cdot \frac{\partial x}{\partial t} \right] \frac{\partial \lambda(s, \tau, t)}{\partial t} ds,
$$
\n(3.7)

in the space  $S^{\theta,\omega}_{\delta,\gamma}$ . Here in the Jacobi matrices  $\frac{\partial x}{\partial t}$ ,  $\partial x$  $rac{\partial x}{\partial \sigma}$ , ∂σ  $\partial \sigma_k$  $\frac{\partial f_{\theta}}{\partial x}$  and partial derivatives  $\frac{\partial f_{\theta}}{\partial \xi_k}$  $\frac{\partial f_{\theta}}{\partial}$  $\partial \eta_k$  $\frac{\partial f_{\theta}}{\partial t}$  $\partial t_j$ in square brackets the variables  $\tau, t, \zeta$  are respectively replaced by  $s, \lambda(s, \tau, t), \mu(s, \tau, \zeta) = (\mu_1, \ldots, \mu_l),$ where  $\mu_k = \mu_k(s, \tau, \sigma_k), \sigma_k = (\xi_k, \eta_k), k = \overline{1, l};$  the vector-function  $f_\theta$  is constructed similarly to formula  $(2.23)$  based on the given vector-function  $f$ .

Note that  $P'_k x, P''_k x, Q_j x$  and  $Rx$  are obtained from  $Tx$  by differentiating it with respect to  $\xi_k, \eta_k, t_j$ and  $\tau$ , respectively.

Integral operators  $(3.2)$ – $(3.7)$  depend on the scalar argument  $\tau$  and the vector arguments  $\zeta_k = (\xi_k, \eta_k), \zeta = (\zeta_1, \ldots, \zeta_l) = ((\xi_1, \eta_1), \ldots, (\xi_l, \eta_l)), t = (t_1, \ldots, t_m),$  $x = (x_1, \ldots, x_n), p = (p_1, \ldots, p_k), p_k =$  $\partial x$  $\partial \xi_k$  $, q = (q_1, \ldots, q_k), q_k =$  $\partial x$  $\partial \eta_k$  $, k = 1, l,$  $r = (r_1, \ldots, r_j), r_j =$  $\partial x$  $\partial t_j$  $i, j = \overline{1,m}$ . Moreover, dependence on  $p, q, r$  is linear.

We have following estimates for  $\lambda(s, \tau, t)$ ,  $\partial \lambda(s,\tau,t)$  $\frac{\partial}{\partial \tau}$ ,  $\partial \lambda(s,\tau,t)$  $\partial t_j$ 

$$
|\lambda(s,\tau,t) - t| \le ||a||_{\circ}|\tau - s|,
$$
  
\n
$$
\left|\frac{\partial\lambda(s,\tau,t)}{\partial\tau}\right| \le ||a||_{\circ} \exp\left[\left\|\frac{\partial a}{\partial t}\right\|_{\circ}|\tau - s|\right],
$$
  
\n
$$
\left|\frac{\partial\lambda(s,\tau,t)}{\partial t_{j}}\right| \le \exp\left[\left\|\frac{\partial a}{\partial t}\right\|_{\circ}|\tau - s|\right], j = \overline{1,m}
$$
\n(3.8)

from the equations of characteristics of the operator D. They are based on the theorems on differentiability with respect to the initial data.

Also we have similar estimates for  $\mu_k(s, \tau, \zeta_k)$ 

$$
|\mu_k(s, \tau, \zeta_k) - z_k(s)| \leq |Z_k(s - \tau)| \cdot |\zeta_k - z_k(\tau)|,
$$
  
\n
$$
\left| \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \tau} \right| \leq |Z_k(s - \tau)| \cdot |- \nu_k I_2 \zeta_k - g_k(\tau)|,
$$
  
\n
$$
\left| \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \xi_k} \right| \leq |Z_k(s - \tau)| \cdot |e'|,
$$
  
\n
$$
\left| \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \eta_k} \right| \leq |Z_k(s - \tau)| \cdot |e''|,
$$
\n(3.9)

where  $|e'| = |e''| = 1$ .

Assume that  $\mu(s, \tau, \zeta) = (\mu_1(s, \tau, \zeta_1), \dots, \mu_l(s, \tau, \zeta_l)).$  Suppose that  $\mu_k = (\mu'_k, \mu''_k)$  and  $z_k =$  $(z'_k, z''_k)$ , then we have the following coordinate representations of the form

$$
\begin{pmatrix}\n\mu'_{k}(s,\tau,\xi_{k},\eta_{k}) - z'_{k}(s) \\
\mu''_{k}(s,\tau,\xi_{k},\eta_{k}) - z''_{k}(s)\n\end{pmatrix} = \begin{pmatrix}\n[\xi_{k} - z'_{k}(\tau)]\cos\nu_{k}(s-\tau) - [\eta_{k} - z''_{k}(\tau)]\sin\nu_{k}(s-\tau) \\
[\xi_{k} - z'_{k}(\tau)]\sin\nu_{k}(s-\tau) + [\eta_{k} - z''_{k}(\tau)]\cos\nu_{k}(s-\tau)\n\end{pmatrix},
$$
\n
$$
\begin{pmatrix}\n\frac{\partial}{\partial\tau}\mu'_{k}(s,\tau,\xi_{k},\eta_{k}) \\
\frac{\partial}{\partial\tau}\mu''_{k}(s,\tau,\xi_{k},\eta_{k})\n\end{pmatrix} = \begin{pmatrix}\n[\nu_{k}\eta_{k} - \varphi_{k}(\tau)]\cos\nu_{k}(s-\tau) + [\nu_{k}\xi_{k} + \psi_{k}(\tau)]\sin\nu_{k}(s-\tau) \\
[\nu_{k}\eta_{k} - \varphi_{k}(\tau)]\sin\nu_{k}(s-\tau) - [\nu_{k}\xi_{k} + \psi_{k}(\tau)]\cos\nu_{k}(s-\tau)\n\end{pmatrix},
$$
\n
$$
\begin{pmatrix}\n\frac{\partial}{\partial\xi_{k}}\mu'_{k}(s,\tau,\xi_{k},\eta_{k}) \\
\frac{\partial}{\partial\xi_{k}}\mu''_{k}(s,\tau,\xi_{k},\eta_{k})\n\end{pmatrix} = \begin{pmatrix}\n\cos\nu_{k}(s-\tau) \\
\sin\nu_{k}(s-\tau)\n\end{pmatrix},
$$
\n
$$
\begin{pmatrix}\n\frac{\partial}{\partial\eta_{k}}\mu'_{k}(s,\tau,\xi_{k},\eta_{k})\n\end{pmatrix} = \begin{pmatrix}\n-\sin\nu_{k}(s-\tau) \\
\cos\nu_{k}(s-\tau)\n\end{pmatrix}.
$$
\n
$$
\begin{pmatrix}\n\frac{\partial}{\partial\eta_{k}}\mu''_{k}(s,\tau,\xi_{k},\eta_{k})\n\end{pmatrix} = \begin{pmatrix}\n-\sin\nu_{k}(s-\tau) \\
\cos\nu_{k}(s-\tau)\n\end{pmatrix}.
$$
\n
$$
\begin{bmatrix}\n\frac{\partial}{\partial\eta_{k}}\mu''_{k}(s,\tau,\xi_{k},\eta_{
$$

Further, we consider the matrix  $\frac{\partial x}{\partial \zeta}$  $\frac{\partial x}{\partial \zeta}$  =  $\partial \xi_k$  $\frac{\partial x_j}{\partial x_j}$  $\partial \eta_k$  $j=1,n$  $_{k=1,l}$ =  $\partial \xi_k$ ,  $\partial \eta_k$  $_{k=1,l}$ whose columns are vectors  $p_k =$  $\partial x$  $\partial \xi_k$ ,  $q_k =$  $\partial x$  $\partial \eta_k$  $, k = 1, l.$  Moreover,

$$
\left|\frac{\partial x}{\partial \zeta}\right| \leq \sum_{k=1}^l \left(\left|\frac{\partial x}{\partial \xi_k}\right| + \left|\frac{\partial x}{\partial \eta_k}\right|\right) = \sum_{k=1}^l (|p_k| + |q_k|).
$$

Similarly, we have

$$
\left|\frac{\partial x}{\partial t}\right| \le \sum_{j=1}^m \left|\frac{\partial x}{\partial t_j}\right| = \sum_{j=1}^m |r_j|.
$$

The matrix 
$$
\frac{\partial \zeta}{\partial \zeta_k} = \left[ \frac{\partial \zeta_j}{\partial \xi_k}, \frac{\partial \zeta_j}{\partial \eta_k} \right]_{j=\overline{1,l}}
$$
 is a  $(l \times 2)$ -dimensional matrix, where the *k*-th row consists of

the units, the rest of its elements are zeros, since  $\zeta_j = (\xi_j, \eta_j), j = \overline{1, j}$ . It is obvious that  $\begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \end{array} \\ \end{array} \\ \begin{array}{c} \end{array} \end{array} \end{array}$ ∂ζ  $\partial \zeta_k$  ≤ √ 2. It is easy to obtain the estimates

$$
\left| \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \xi_k} \right| \leq \sqrt{2},
$$
\n
$$
\left| \frac{\partial \mu_k(s, \tau, \zeta_k)}{\partial \eta_k} \right| \leq \sqrt{2},
$$
\n
$$
\left| \frac{\partial \mu(s, \tau, \zeta)}{\partial \tau} \right| \leq \sqrt{2l}, k = \overline{1, l}
$$
\n(3.10)

from the above representations and estimates (3.9).

11)

 $\partial x$ 

 $\partial x$ 

∂ζ

 $k=1$ 

 $(|p_k| + |q_k|)$ 

Further, by virtue of condition (3.1) on twice continuous differentiability, we have the following estimates for the differences and their corollaries

1) 
$$
|f_{\theta}(\tau, t, \zeta, x) - f_{\theta}(\tau, t, \zeta, y)| \leq c_1 |x - y|
$$
,  
\n2)  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial \zeta} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial t} \right| \leq c_2 |x - y|$ ,  
\n3)  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial t_{j}} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial t_{j}} \right| \leq c_3 |x - y|, j = \overline{1, m}$ ,  
\n4)  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \right| \leq c_4 |x - y|$ ,  
\n5)  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} \frac{\partial x}{\partial \zeta} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \frac{\partial y}{\partial \zeta} \right| \leq \left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \right| \cdot \left| \frac{\partial x}{\partial \zeta} \right|$   
\n+  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \frac{\partial x}{\partial t} - \frac{\partial y}{\partial x} \right|$ ,  
\n6)  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} \frac{\partial x}{\partial t} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \frac{\partial y}{\partial t} \right| \leq \left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \right| \cdot \left| \frac{\partial x}{\partial t} \right|$   
\n+  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} \right| \cdot \left| \frac{\partial x}{\partial t} - \frac{\partial y}{\partial t} \right|$ ,  
\n7) 

+ 
$$
(\chi_4 + c_4\gamma) \sum_{k=1}^{l} (|p_k - \tilde{p}_k| + |q_k - \tilde{q}_k|)
$$
, where  $\tilde{p}_k = \frac{\partial y}{\partial \xi_k}$ ,  $\tilde{q}_k = \frac{\partial y}{\partial \eta_k}$ ,  $k = \overline{1, l}$ ,  
\n12)  $\left| \frac{\partial f_{\theta}(\tau, t, \zeta, x)}{\partial x} \frac{\partial x}{\partial t} - \frac{\partial f_{\theta}(\tau, t, \zeta, y)}{\partial x} \frac{\partial y}{\partial t} \right| \le c_4 |x - y| \sum_{j=1}^{m} |r_j| + (\chi_4 + c_4\gamma) \sum_{j=1}^{m} |r_j - \tilde{r}_j|$ ,  
\nwhere  $\tilde{r}_j = \frac{\partial y}{\partial t_j}$ ,  $j = \overline{1, m}$ . Here  $c_1, c_2, c_3, c_4$  are the positive Lipschitz constants of  
\nthe vector and matrix functions  $f_{\theta}$ ,  $\frac{\partial f_{\theta}}{\partial \zeta}$ ,  $\frac{\partial f_{\theta}}{\partial t_j}$ ,  $j = \overline{1, m}$ ,  $\frac{\partial f_{\theta}}{\partial x}$ ;  $\chi_1, \chi_2, \chi_3, \chi_4 > 0$  are the maximums of  
\nthe functions  $|f_{\theta}|$ ,  $\left| \frac{\partial f_{\theta}}{\partial \zeta} \right|$ ,  $\left| \frac{\partial f_{\theta}}{\partial t_j} \right|$ ,  $j = \overline{1, m}$ ,  $\left| \frac{\partial f_{\theta}}{\partial x} \right|$  with respect to  $(\tau, t, \zeta, 0)$ . (For  $\chi_3$  the maximum is  
\nalso with respect to  $j = 1, m$ .)

Obviously, the integral operator  $T_\circ$  of the form

$$
(T_{\circ}x)(\tau, t, \zeta) = \int\limits_{\tau}^{\tau+\theta} K(\tau, s)x(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta)) ds \qquad (3.11)
$$

with kernel (3.3) is linear, maps the space  $S^{\theta,\omega}_{\delta,\gamma}$  into itself and is bounded by

$$
||T_{\circ}x||_{\circ} = \kappa ||x||_{\circ},\tag{3.12}
$$

where

$$
\kappa = \sup_{0 \le \tau \le \theta} \int\limits_{\tau}^{\tau+\theta} |K(\tau, s)| \, ds.
$$

Now we estimate operators (3.2)–(3.7) by using inequalities 1)-12) and (3.8)–(3.12). For  $x \in S^{\theta,\omega}_{\delta,\gamma}$  $_{\delta,\gamma}$  $|(Tx)(\tau, t, \zeta)| \leq \kappa (\chi_1 + c_1 ||x||_{\circ}),$ 

$$
|(P'_{k}x)(\tau,t,\zeta)| \leq \kappa \Big[\chi_{2}+c_{2}\|x\|_{\diamond}+\left(\chi_{4}+c_{4}\|x\|_{\diamond}\right)\sum_{k=1}^{l}\left(\|p_{k}\|_{\diamond}-\|q_{k}\|_{\diamond}\right)\Big]\cdot\sqrt{2}\cdot\sqrt{2},
$$
  

$$
|(P''_{k}x)(\tau,t,\zeta)| \leq \kappa \Big[\chi_{2}+c_{2}\|x\|_{\diamond}+\left(\chi_{4}+c_{4}\|x\|_{\diamond}\right)\sum_{k=1}^{l}\left(\|p_{k}\|_{\diamond}-\|q_{k}\|_{\diamond}\right)\Big]\cdot\sqrt{2}\cdot\sqrt{2},
$$
  

$$
|(Q_{j}x)(\tau,t,\zeta)| \leq \kappa \Big[\chi_{3}+c_{3}\|x\|_{\diamond}+\left(\chi_{4}+c_{4}\|x\|_{\diamond}\right)\sum_{j=1}^{m}\|r_{j}\|_{\diamond}\Big]\cdot e^{2\left\|\frac{\partial a}{\partial t}\right\|_{\diamond}\theta},
$$
  

$$
|(Rx)(\tau,t,\zeta)| \leq \kappa \Big[\chi_{2}+c_{2}\|x\|_{\diamond}+\left(\chi_{4}+c_{4}\|x\|_{\diamond}\right)\sum_{k=1}^{l}\left(\|p_{k}\|_{\diamond}-\|q_{k}\|_{\diamond}\right)\Big]\cdot\sqrt{2l}
$$
  

$$
+\kappa \Big[\chi_{3}+c_{3}\|x\|_{\diamond}+\left(\chi_{4}+c_{4}\|x\|_{\diamond}\right)\sum_{j=1}^{m}\|r_{j}\|_{\diamond}\Big]\cdot e^{2\left\|\frac{\partial a}{\partial t}\right\|_{\diamond}\theta}+|A|\kappa\big(\chi_{1}+c_{1}\|x\|_{\diamond}\big)
$$
  

$$
+\left(\chi_{1}+c_{1}\|x\|_{\diamond}\right) \text{ for } |s| \leq \theta, |\tau| \leq \theta.
$$
  
Let  $\max\{\chi_{1},\chi_{2},\chi_{3},\chi_{4}\}=\chi, \max\{c_{1},c_{2},c_{3},c_{4}\}=c.$  Then we obtain the follow

wing estimates via γ:

$$
||(Tx)||_{\circ} \le \kappa(\chi + c\gamma)(1 + \gamma);
$$
  

$$
||(P'_kx)||_{\circ} \le 4l\kappa(\chi + c\gamma)(1 + \gamma), k = \overline{1,l};
$$

$$
||(P_k''x)||_{\circ} \le 4l\kappa(\chi + c\gamma)(1+\gamma), k = \overline{1,l};
$$
  

$$
||(Q_jx)||_{\circ} \le m\varepsilon\kappa(\chi + c\gamma)(1+\gamma), j = \overline{1,m}, \varepsilon = exp\left[2\left\|\frac{\partial a}{\partial t}\right\|_{\circ}\theta\right];
$$
  

$$
||(Rx)||_{\circ} \le [1 + \kappa|A| + m\varepsilon\kappa(1+\gamma) + (2l)^{3/2}\kappa(1+\gamma)](\chi + c\gamma)
$$

 $\leq (1 + \kappa)A + m\varepsilon\kappa + (2l)^{3/2}\kappa(\chi + c\gamma)(1 + \gamma).$ Then we have

$$
||Tx|| = ||(Tx)||_{\circ} + \sum_{k=1}^{l} (||(P'_{k}x)||_{\circ} + ||(P''_{k}x)||_{\circ}) + \sum_{j=1}^{m} ||(Q_{j}x)||_{\circ}
$$
  
+ 
$$
||(Rx)||_{\circ} \leq (\kappa + 8l^{2}\kappa + m^{2}\varepsilon\kappa + 1 + \kappa|A| + m\varepsilon\kappa
$$
  
+ 
$$
(3.13)
$$
  
+ 
$$
||(x||_{\circ})^{3/2}\kappa(1+\gamma)(\chi + c\gamma) = c_{*}(\chi + c\gamma) \leq \gamma,
$$

for sufficiently small  $\chi$  and  $c$ , where

$$
c_* = (\kappa + 8l^2\kappa + m^2\varepsilon\kappa + 1 + \kappa|A| + m\varepsilon\kappa + (2l)^{3/2}\kappa(1+\gamma).
$$

So, for sufficiently small  $\chi$  and c the operator T maps the space  $S^{\theta,\omega}_{\delta,\gamma}$  into itself.

Next, we estimate the differences in the values of operators  $(3.2)$ – $(3.7)$  at different points  $x, y \in$  $S^{\theta,\omega}_{\delta,\gamma}$  by using 1)-12) and (3.8)-(3.12):

$$
|(Tx)(\tau, t, \zeta) - (Ty)(\tau, t, \zeta)| \leq \kappa c_1 ||x - y||_0,
$$
  
\n
$$
|(P'_kx)(\tau, t, \zeta) - (P'_ky)(\tau, t, \zeta)| \leq \kappa [c_2 ||x - y||_0 + c_4 ||x - y||_0 \sum_{k=1}^l (||p_k||_0 + ||q_k||_0)
$$
  
\n
$$
+ (\chi_4 + c_4 ||y||_0) \cdot \sum_{k=1}^l (||p_k - \tilde{p}_k||_0 + ||q_k - \tilde{q}_k||_0) \cdot \sqrt{2} \cdot \sqrt{2},
$$
  
\n
$$
|(P''_kx)(\tau, t, \zeta) - (P''_ky)(\tau, t, \zeta)| \leq \kappa [c_2 ||x - y||_0 + c_4 ||x - y||_0 \sum_{k=1}^l (||p_k||_0 + ||q_k||_0)
$$
  
\n
$$
+ (\chi_4 + c_4 ||y||_0) \cdot \sum_{k=1}^l (||p_k - \tilde{p}_k||_0 + ||q_k - \tilde{q}_k||_0) \cdot \sqrt{2} \cdot \sqrt{2},
$$
  
\n
$$
|(Q_jx)(\tau, t, \zeta) - (Q_jy)(\tau, t, \zeta)| \leq \kappa [c_3 ||x - y||_0 + c_4 ||x - y||_0 \sum_{j=1}^m ||r_j||_0
$$
  
\n
$$
+ (\chi_4 + c_4 ||y||_0) \sum_{j=1}^m ||r_j - \tilde{r}_j||_0 \cdot e^{2||\frac{\partial \alpha}{\partial t}||_{\infty}^{\theta}},
$$
  
\n
$$
|(Rx)(\tau, t, \zeta) - (Ry)(\tau, t, \zeta)| \leq c_1 ||x - y||_0 + \kappa |A| \cdot c_1 ||x - y||_0 + \kappa [c_2 ||x - y||_0
$$
  
\n
$$
+ c_4 ||x - y||_0 \sum_{k=1}^l (||p_k||_0 + ||q_k||_0) + (\chi_4 + c_4 ||y||_0) \sum_{k=1}^l (||p_k - \tilde{p}_k||_0 + ||q_k - \tilde{q}_k||_0) \cdot \sqrt{2l}
$$
  
\n

 $\|Tx - Ty\|_{\circ} \leq \kappa c \|x - y\|_{\circ};$ 

$$
||P_{k}^{r}x - P_{k}^{r}y||_{\infty} \leq 2\kappa \Big[(c + c2l\gamma)||x - y||_{\infty} + (\chi + c\gamma)\sum_{k=1}^{l} (||p_{k} - \tilde{p}_{k}||_{\infty} + ||q_{k} - \tilde{q}_{k}||_{\infty})\Big],
$$
  
\n
$$
||P_{k}^{r}x - P_{k}^{r}y||_{\infty} \leq 2\kappa \Big[(c + c2l\gamma)||x - y||_{\infty} + (\chi + c\gamma)\sum_{k=1}^{l} (||p_{k} - \tilde{p}_{k}||_{\infty} + ||q_{k} - \tilde{q}_{k}||_{\infty})\Big],
$$
  
\n
$$
k = \overline{1, l};
$$
  
\n
$$
||Q_{j}x - Q_{j}y||_{\infty} \leq \kappa \varepsilon \Big[(c + c m\gamma)||x - y||_{\infty} + (\chi + c\gamma)\sum_{j=1}^{m} ||r_{j} - \tilde{r}_{j}||_{\infty}\Big], j = \overline{1, m};
$$
  
\n
$$
||Rx - Ry||_{\infty} \leq \Big[c + \kappa c|A| + \kappa c\sqrt{2l} + \kappa c(2l)^{3/2}\gamma)\Big] ||x - y||_{\infty}
$$
  
\n
$$
+ \kappa \sqrt{2l}(\chi + c\gamma)\sum_{j=1}^{l} (||p_{k} - \tilde{p}_{k}||_{\infty} + ||q_{k} - \tilde{q}_{k}||_{\infty}) + \kappa ||a||_{\infty}(\varepsilon + c m\gamma)||x - y||_{\infty}
$$
  
\n
$$
+ \kappa ||a||_{\infty}(\chi + c\gamma)\sum_{j=1}^{m} ||r_{j} - \tilde{r}_{j}||_{\infty}.
$$
  
\nFurther, summing up these inequalities, we obtain  
\n
$$
||Tx - Ty||_{\infty} + \sum_{k=1}^{l} ||P_{k}^{r}x - P_{k}^{r}y||_{\infty} + \sum_{k=1}^{l} ||P_{k}^{r}x - P_{k}^{r}y||_{\infty} + \sum_{j=1}^{m} ||Q_{j}x - Q_{j}y||_{\infty} + ||Rx - Ry||_{\infty}
$$

Hence, for sufficiently small values of  $\chi$  and  $c$ , we have

$$
||Tx - Ty|| = ||Tx - Ty||_{\circ} + \sum_{k=1}^{l} ||P'_{k}x - P'_{k}y||_{\circ} + \sum_{k=1}^{l} ||P''_{k}x - P''_{k}y||_{\circ}
$$
  
+ 
$$
\sum_{j=1}^{m} ||Q_{j}x - Q_{j}y||_{\circ} + ||Rx - Ry||_{\circ} \leq (\chi + c\gamma)c_{0} [||x - y||_{\circ}
$$
  
+ 
$$
\sum_{k=1}^{l} (||p_{k} - \tilde{p}_{k}||_{\circ} + ||q_{k} - \tilde{q}_{k}||_{\circ}) + \sum_{j=1}^{m} ||r_{j} - \tilde{r}_{j}||_{\circ} \leq (\chi + c\gamma)c_{0} ||x - y||,
$$
  
(3.14)

where  $c_0 = max\Big\{ \big[ \kappa + 8l\kappa(1 + 2l\gamma) + \kappa \varepsilon(1 + m\gamma) + (1 + \kappa|A| + \kappa) \Big\}$ √  $\overline{2l}$  +  $\kappa(2l)^{3/2}\gamma$ )  $+ \kappa ||a||_{\circ} \varepsilon (1 + m\gamma)$ ; [8 $l\kappa +$ √  $\overline{2l}\kappa$ ];  $\lceil \kappa \varepsilon m + \kappa m \varepsilon ||a||_{\circ} \rceil$   $\}$ .

Under the conditions

$$
c_*(\chi + c\gamma) \le \gamma, c_0(\chi + c\gamma) < 1\tag{3.15}
$$

by virtue of estimates (3.13) and (3.14), it is clear that operator (3.2) maps the space  $S_{\delta,\gamma}^{\theta,\omega}$  into itself and is a contraction operator. The space  $S_{\delta,\gamma}^{\theta,\omega}$  is complete. Therefore, there exists a unique fixed point  $x^*(\tau,t,\zeta) \in S$  $\frac{\theta,\omega}{\delta,\gamma}$  of the operator T:  $(Tx^*)(\tau, t, \zeta) = x^*(\tau, t, \zeta)$  that is

$$
x^*(\tau, t, \zeta) \equiv \int_{\tau}^{\tau+\theta} K(\tau, s) f_{\theta}(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta), x^*(s, \lambda(s, \tau, t), \mu(s, \tau, \zeta))) ds.
$$

It is the unique  $(\theta, \omega)$ -periodic solution of system  $(2.1)$  by condition 3<sup>0</sup>. Thus, the following theorem is proved.

#### Theorem 3.1. Suppose that

1) conditions  $(2.3)-(2.6)$  are satisfied with respect to the input data of the differentiation operator  $D,$ 

2) the matrix A is such that linear homogeneous system  $(2.15)$  corresponding to the given quasilinear system has no  $(\theta, \omega)$ -periodic solutions, except for the zero one (in particular, if conditions  $(2.18)$ – $(2.20)$  are satisfied)

and

3) the vector-function  $f(\tau, t, \zeta, x)$  satisfies condition (3.1).

Then, under condition (3.15) with respect to the parameters  $c_0$ ,  $c_*, c, \chi, \gamma$  quasilinear system (2.1) has a unique  $(\theta, \omega)$ -periodic solution  $x^*(\tau, t, \zeta) \in S^{\theta, \omega}_{\delta, \gamma}$ .

#### 4 Conclusion

A technique has been developed for investigation of oscillatory solutions of perturbed quasilinear autonomous systems of form  $(2.1)$ – $(2.2)$ , based on proving the existence of a fixed point of nonlinear operator (3.2), which is an analogue of representation (2.24) in space  $S^{\theta,\omega}_{\delta,\gamma}$ .

The main essence of the technique for investigation of multiperiodic solutions of the system under consideration is the further development of the methods of papers [15, 16], on the basis of the fundamental methods of works [4, 9, 11, 17] and investigation carried out in [5–8, 10, 12–14]. In conclusion, the developed technique allows to establish sufficient conditions for the existence and uniqueness of multiperiodic solutions of quasilinear systems (2.1) with the differentiation operator  $(2.2).$ 

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