

ON DIRECT VARIATIONAL FORMULATIONS  
FOR SECOND ORDER EVOLUTIONARY EQUATIONS

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**Abstract.** The existence of direct variational formulations for a wide class of second order evolutionary equations is investigated.

## 1 Introduction

Consider the following operator equation

$$N(u) \equiv P_{2u,t}u_{tt} + P_{1u,t}u_t + Q(t, u) = 0_V, \quad (1.1)$$

$$u \in D(N) \subseteq U \subseteq V, \quad t \in [t_0, t_1] \subset \mathbb{R}; \quad u_t \equiv D_t u \equiv \frac{d}{dt}u, \quad u_{tt} \equiv \frac{d^2}{dt^2}u.$$

Here  $\forall t \in [t_0, t_1], \forall u \in U_1$   $P_{iu,t} : U_1 \rightarrow V_1$  ( $i = 1, 2$ ) are linear operators;  $Q : [t_0, t_1] \times U_1 \rightarrow V_1$  is an arbitrary operator;  $D(N)$  is a domain of definition of the operator  $N : D(N) \subseteq U \rightarrow V$ ;  $U = C^2([t_0, t_1]; U_1)$ ,  $V = C([t_0, t_1]; V_1)$ ,  $U_1, V_1$  are linear normed spaces,  $U_1 \subseteq V_1$ .

Assume that for every  $t \in (t_0, t_1)$  and  $g(t), u(t) \in U_1$  the function  $P_{1u,t}g(t)$  is continuously differentiable and  $P_{2u,t}g(t)$  is twice continuously differentiable on  $(t_0, t_1)$ .

Any function  $u \in D(N)$  is called a solution of problem (1.1) if it satisfies equation (1.1).

In the sequel, we shall write

$$N(u) \equiv P_{2u}u_{tt} + P_{1u}u_t + Q(u) = 0_V,$$

bearing in mind that the operators  $P_{1u}, P_{2u}$  and  $Q$  may also depend on  $t$ .

First let us introduce the following concepts. Let  $N$  be an operator such that its domain of definition  $D(N) \subseteq U$  and the range of values  $R(N) \subseteq V$ , where  $U$  and  $V$  are linear normed spaces over  $\mathbb{R}$ , i. e.

$$N(u) = v, \quad u \in U, \quad v \in V.$$

If there exists the limit

$$\delta N(u, h) = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \{N(u + \varepsilon h) - N(u)\}, \quad u \in D(N), \quad (u + \varepsilon h) \in D(N), \quad (1.2)$$

then it is called the Gâteaux variation of the operator  $N$  at the point  $u$  or the first variation of the operator  $N$  at the point  $u$ .

$\delta N(u, h)$  is homogeneous with respect to  $h$  :  $\delta N(u, \lambda h) = \lambda \delta N(u, h)$ , but the operator  $\delta N(u, \cdot) : U \rightarrow V$  is not always additive with respect to  $h$ .

If  $\delta N(u, h)$  is a linear operator with respect to  $h$ , when  $u$  is a fixed element of  $D(N)$ , then we say that the operator  $N$  is Gâteaux differentiable at the point  $u$ . The expression  $\delta N(u, h)$  is called the Gâteaux differential and denoted by  $DN(u, h)$ . In this case we shall also write  $DN(u, h) = N'_u h$  and say that  $N'_u$  is the Gâteaux derivative of operator  $N$  at the point  $u$ .

If  $N$  is a linear operator then  $N'_u h = Nh$ , i. e. the Gâteaux derivative of the linear operator coincides with it.

Further assume that for any given operator  $N : D(N) \subset U \rightarrow V$  there exists its Gâteaux derivative at any point  $u \in D(N)$ . The domain of definition  $D(N'_u)$  consists of elements  $h \in U$  such that  $(u + \varepsilon h) \in D(N)$  for all  $\varepsilon$  sufficiently small. In this case  $h \in D(N'_u)$  is called an admissible element.

If the Gâteaux derivative of the operator  $N$  exists, then the following equality holds

$$N(u + \varepsilon h) = N(u) + \varepsilon N'_u h + r(u, \varepsilon h), \quad u \in D(N), \quad (1.3)$$

where for any fixed element  $h \in D(N'_u)$

$$\lim_{\varepsilon \rightarrow 0} \frac{r(u, \varepsilon h)}{\varepsilon} = 0_V.$$

If the Gâteaux derivative of the operator  $N$  is known and  $0_U \in D(N)$ , then

$$N(u) = \int_0^1 N'_{tu} u dt + N(0_U). \quad (1.4)$$

Note that for any linear operator  $\tilde{N}_u$  which may depend on  $u$  in a nonlinear way, the Gâteaux derivative is defined by

$$\tilde{N}'_u(g; h) = \lim_{\varepsilon \rightarrow 0} \frac{\tilde{N}_{u+\varepsilon h} g - \tilde{N}_u g}{\varepsilon}. \quad (1.5)$$

The second Gâteaux derivative  $N''_u$  of the operator  $N$  is given by

$$N''_u(h_1, h_2) = \frac{\partial^2}{\partial \varepsilon^1 \partial \varepsilon^2} N(u + \varepsilon^1 h_1 + \varepsilon^2 h_2) |_{\varepsilon^1 = \varepsilon^2 = 0}. \quad (1.6)$$

In the most of applications  $N''_u$  satisfies the symmetry condition

$$N''_u(h_1, h_2) = N''_u(h_2, h_1).$$

Now we need some notation and notions for bilinear forms and potential operators.

**Definition 1.** A mapping  $\Phi(\cdot, \cdot) : V \times U \rightarrow \mathbb{R}$  is said to be a bilinear form if it is linear with respect to every argument.

**Definition 2.** A bilinear form  $\Phi(\cdot, \cdot) : V \times V \rightarrow \mathbb{R}$  is called symmetric if

$$\Phi(v, g) = \Phi(g, v) \quad \forall g, v \in V.$$

Consider a bilinear form

$$\Phi(\cdot, \cdot) \equiv \int_{t_0}^{t_1} \langle \cdot, \cdot \rangle dt : V \times U \rightarrow \mathbb{R} \quad (1.7)$$

such that the bilinear mapping  $\Phi_1(\cdot, \cdot) \equiv \langle \cdot, \cdot \rangle$  satisfies the following conditions:

$$\langle v_1(t), v_2(t) \rangle = \langle v_2(t), v_1(t) \rangle \quad \forall v_1(t), v_2(t) \in V_1, \quad (1.8)$$

$$D_t \langle v(t), g(t) \rangle = \langle D_t v(t), g(t) \rangle + \langle v(t), D_t g(t) \rangle \quad \forall v, g \in C^1([t_0, t_1]; U_1). \quad (1.9)$$

If  $v = v(x, t), x \in \Omega \subset \mathbb{R}^n, t \in (t_0, t_1), U_1 = V_1 = C(\overline{\Omega})$ , then we can take for example

$$\langle v, g \rangle = \int_{\Omega} v(x, t)g(x, t) dx. \quad (1.10)$$

**Definition 3.** The operator  $N : D(N) \subset U \rightarrow V$  is said to be potential on the set  $D(N)$  with respect to the bilinear form  $\Phi(\cdot, \cdot) : V \times U \rightarrow \mathbb{R}$ , if there exists a functional  $F_N : D(F_N) = D(N) \rightarrow \mathbb{R}$  such that

$$\delta F_N[u, h] = \Phi(N(u), h) \quad \forall u \in D(N), \quad \forall h \in D(N'_u).$$

The functional  $F_N$  is called the potential of the operator  $N$ , and in its turn the operator  $N$  is called the gradient of the functional  $F_N$ . In this case we shall write  $N = \text{grad}_{\Phi} F_N$ .

An element  $u \in D(N)$  such that  $\delta F_N[u, h] = 0 \quad \forall h \in D(N'_u)$ , is said to be a critical point of the functional  $F_N$ .

The following theorem is needed in the sequel.

**Theorem 1.1. [3, 5]** Consider the operator  $N : D(N) \subset U \rightarrow V$  and the bilinear form  $\Phi(\cdot, \cdot) : V \times U \rightarrow \mathbb{R}$  such that for any fixed elements  $u \in D(N), g, h \in D(N'_u)$  the function  $\psi(\varepsilon) = \Phi(N(u + \varepsilon h), g)$  belongs to the class  $C^1[0, 1]$ . For  $N$  to be potential on the convex open set  $D(N)$  with respect to  $\Phi$  it is necessary and sufficient that

$$\Phi(N'_u h, g) = \Phi(N'_u g, h) \quad \forall u \in D(N), \quad \forall g, h \in D(N'_u). \quad (1.11)$$

Under this condition the potential  $F_N$  is given by

$$F_N[u] = \int_0^1 \Phi(N(u_0 + \lambda(u - u_0)), u - u_0) d\lambda + F_N[u_0], \quad (1.12)$$

where  $u_0$  is a fixed element of  $D(N)$ .

## 2 Main results

Denoting by  $(\dots)^*$  the operator adjoint to the operator  $(\dots)$ , we shall prove

**Theorem 2.1.** *Suppose that  $D_t^* = -D_t$  on  $D(N'_u)$ ; then for operator (1.1) to be potential on  $D(N)$  with respect to bilinear form (1.7) it is necessary and sufficient that on  $D(N'_u)$*

$$P_{2u} - P_{2u}^* = 0, \quad (2.1)$$

$$P_{2u}^{*'}(\cdot; u_t) = 0, \quad (2.2)$$

$$-2\frac{\partial P_{2u}^*}{\partial t} + P_{1u}^* + P_{1u} = 0, \quad (2.3)$$

$$-\frac{\partial^2 P_{2u}^*}{\partial t^2} + \frac{\partial P_{1u}^*}{\partial t} + Q'_u - Q_u^* = 0, \quad (2.4)$$

$$-\left(\frac{\partial P_{2u}^*}{\partial t}\right)'_u(\cdot; u_t) - \frac{\partial P_{2u}^{*'}}{\partial t}(\cdot; u_t) + P_{1u}^{*'}(\cdot; u_t) + P'_{1u}(u_t; \cdot) - [P'_{1u}(u_t; \cdot)]^* = 0, \quad (2.5)$$

$$P'_{2u}(u_{tt}; \cdot) - P_{2u}^{*'}(\cdot; u_{tt}) - [P'_{2u}(u_{tt}; \cdot)]^* = 0. \quad (2.6)$$

$$\forall u \in D(N), \quad \forall t \in [t_0, t_1].$$

*Proof.* Using (1.1) and (1.5), we get

$$N'_u h = P_{2u} h_{tt} + P'_{2u}(u_{tt}; h) + P_{1u} h_t + P'_{1u}(u_t; h) + Q'_u h.$$

In this case (1.11) can be written in the form

$$\begin{aligned} & \int_{t_0}^{t_1} \langle P_{2u} h_{tt} + P'_{2u}(u_{tt}; h) + P_{1u} h_t + P'_{1u}(u_t; h) + Q'_u h, g \rangle dt \\ &= \int_{t_0}^{t_1} \langle P_{2u} g_{tt} + P'_{2u}(u_{tt}; g) + P_{1u} g_t + P'_{1u}(u_t; g) + Q'_u g, h \rangle dt, \end{aligned}$$

or

$$\begin{aligned} & \int_{t_0}^{t_1} \{ \langle P_{2u} h_{tt} + P'_{2u}(u_{tt}; h) + P_{1u} h_t + P'_{1u}(u_t; h) + Q'_u h, g \rangle \\ & - \langle D_t^2(P_{2u}^* h) + [P'_{2u}(u_{tt}; \cdot)]^* h - D_t(P_{1u}^* h) + [P'_{1u}(u_t; \cdot)]^* h + Q_u^{*'} h, g \rangle \} dt = 0 \quad (2.7) \\ & \forall u \in D(N), \quad \forall g, h \in D(N'_u). \end{aligned}$$

Taking into account the second Gâteaux derivative, we obtain

$$D_t^2(P_{2u}^* h) = D_t[D_t(P_{2u}^* h)] = D_t \left[ P_{2u}^* h_t + \frac{\partial P_{2u}^*}{\partial t} h + P_{2u}^{*'}(h; u_t) \right]$$

$$\begin{aligned}
&= P_{2u}^* h_{tt} + 2 \frac{\partial P_{2u}^*}{\partial t} h_t + 2 P_{2u}^{*'}(h_t; u_t) + \frac{\partial^2 P_{2u}^*}{\partial t^2} h + \left( \frac{\partial P_{2u}^*}{\partial t} \right)'_u (h; u_t) + P_{2u}^{*''}(h; u_t; u_t) \\
&\quad + P_{2u}^{*'}(h; u_{tt}) + \frac{\partial P_{2u}^{*'}}{\partial t}(h; u_t). \tag{2.8}
\end{aligned}$$

From (2.8), (2.7) we get the following:

$$\begin{aligned}
&\int_{t_0}^{t_1} < P_{2u} h_{tt} + P_{2u}'(u_{tt}; h) + P_{1u} h_t + P_{1u}'(u_t; h) + Q'_u h - P_{2u}^* h_{tt} - 2 \frac{\partial P_{2u}^*}{\partial t} h_t \\
&\quad - 2 P_{2u}^{*'}(h_t; u_t) - \frac{\partial^2 P_{2u}^*}{\partial t^2} h - \left( \frac{\partial P_{2u}^*}{\partial t} \right)'_u (h; u_t) - P_{2u}^{*''}(h; u_t; u_t) - P_{2u}^{*'}(h; u_{tt}) \\
&\quad - \frac{\partial P_{2u}^{*'}}{\partial t}(h; u_t) - [P_{2u}'(u_{tt}; \cdot)]^* h + P_{1u}^* h_t + \frac{\partial P_{1u}^*}{\partial t} h + P_{1u}'(h; u_t) - [P_{1u}'(u_t; \cdot)]^* h \\
&\quad - Q_u'^* h, g > dt = 0.
\end{aligned}$$

Thus condition (2.7) is represented in the form

$$\begin{aligned}
&\int_{t_0}^{t_1} < \{ (P_{2u} - P_{2u}^*) D_{tt} + \left( P_{1u} - 2 \frac{\partial P_{2u}^*}{\partial t} - 2 P_{2u}^{*'}(\cdot; u_t) + P_{1u}^* \right) D_t \\
&\quad + P_{2u}'(u_{tt}; \cdot) + P_{1u}'(u_t; \cdot) + Q'_u - \frac{\partial^2 P_{2u}^*}{\partial t^2} - \left( \frac{\partial P_{2u}^*}{\partial t} \right)'_u (\cdot; u_t) - P_{2u}^{*''}(\cdot; u_t; u_t) \\
&\quad - P_{2u}^{*'}(\cdot; u_{tt}) - \frac{\partial P_{2u}^{*'}}{\partial t}(\cdot; u_t) - [P_{2u}'(u_{tt}; \cdot)]^* + \frac{\partial P_{1u}^*}{\partial t} + P_{1u}'(\cdot; u_t) - [P_{1u}'(u_t; \cdot)]^* \\
&\quad - Q_u'^* \} h, g > dt = 0 \qquad \forall u \in D(N), \quad \forall g, h \in D(N'_u).
\end{aligned}$$

This is fulfilled identically if and only if

$$\begin{aligned}
&\left[ (P_{2u} - P_{2u}^*) D_{tt} + \left( P_{1u} - 2 \frac{\partial P_{2u}^*}{\partial t} - 2 P_{2u}^{*'}(\cdot; u_t) + P_{1u}^* \right) D_t \right. \\
&\quad + P_{2u}'(u_{tt}; \cdot) + P_{1u}'(u_t; \cdot) + Q'_u - \frac{\partial^2 P_{2u}^*}{\partial t^2} - \left( \frac{\partial P_{2u}^*}{\partial t} \right)'_u (\cdot; u_t) - P_{2u}^{*''}(\cdot; u_t; u_t) \\
&\quad - P_{2u}^{*'}(\cdot; u_{tt}) - \frac{\partial P_{2u}^{*'}}{\partial t}(\cdot; u_t) - [P_{2u}'(u_{tt}; \cdot)]^* + \frac{\partial P_{1u}^*}{\partial t} + P_{1u}'(\cdot; u_t) - [P_{1u}'(u_t; \cdot)]^* \\
&\quad \left. - Q_u'^* \right] h = 0_V \qquad \forall u \in D(N), \quad \forall h \in D(N'_u).
\end{aligned}$$

The necessary and sufficient conditions for this equality to be valid are that conditions (2.1) - (2.6) hold.  $\square$

**Remark 1.** By (2.1), (1.4), and (2.2) it follows that

$$P_{2u} h - P_{20} h = \int_0^1 \frac{d}{d\varepsilon} P_{2\varepsilon u} h d\varepsilon = \int_0^1 P_{2\varepsilon u}'(h; u) d\varepsilon = 0,$$

i.e. the operator  $P_{2u}$  does not depend on  $u$  on  $D(N'_u)$ .

Therefore conditions (2.1) - (2.6) are equivalent to the following ones:

$$P_2 - P_2^* = 0, \quad (2.9)$$

$$-2 \frac{\partial P_2}{\partial t} + P_{1u}^* + P_{1u} = 0, \quad (2.10)$$

$$\frac{\partial^2 P_2}{\partial t^2} - \frac{\partial P_{1u}}{\partial t} + Q'_u - Q'^*_u = 0, \quad (2.11)$$

$$-P'_{1u}(\cdot; u_t) + P'_{1u}(u_t; \cdot) - [P'_{1u}(u_t; \cdot)]^* = 0, \quad (2.12)$$

$$P'_{2u}(u_{tt}; \cdot) - [P'_{2u}(u_{tt}; \cdot)]^* = 0. \quad (2.13)$$

**Theorem 2.2.** *If  $D_t^* = -D_t$  on  $D(N'_u)$ , then conditions (2.1) - (2.6) hold if and only if equation (1.1) can be represented in the form*

$$N(u) \equiv P_{2u}u_{tt} + P_{1u}u_t + Q(u) \quad (2.14)$$

$$\equiv (-\mathcal{R}_2 - \mathcal{R}_2^*)u_{tt} + (\mathcal{R}'_{1u} - \mathcal{R}'_{1u} - 2 \frac{\partial \mathcal{R}_2^*}{\partial t})u_t + \text{grad}_{\Phi_1} \mathcal{B}[u] - \frac{\partial \mathcal{R}_1}{\partial t}(u) + \frac{\partial^2 \mathcal{R}_2}{\partial t^2}u = 0_V$$

$$\forall u \in D(N).$$

The operators  $\mathcal{R}_1, \mathcal{R}_2, \mathcal{B}$  are defined by

$$\Phi(\mathcal{R}_1(u), u_t) = \int_{t_0}^{t_1} \int_0^1 \left\langle -P_{1\tilde{u}(\lambda)}(u - u_0), \frac{\partial \tilde{u}(\lambda)}{\partial t} \right\rangle d\lambda dt, \quad (2.15)$$

$$\Phi(\mathcal{R}_2 u_t, u_t) = \int_{t_0}^{t_1} \int_0^1 \left\langle -P_2(u_t - u_{0t}), \frac{\partial \tilde{u}(\lambda)}{\partial t} \right\rangle d\lambda dt, \quad (2.16)$$

$$\mathcal{B}[u] = \int_0^1 \left\langle Q(\tilde{u}(\lambda)) + \lambda \frac{\partial P_{1\tilde{u}(\lambda)}}{\partial t}(u - u_0) - \lambda \frac{\partial^2 P_2}{\partial t^2}(u - u_0), u - u_0 \right\rangle d\lambda, \quad (2.17)$$

where  $\tilde{u}(\lambda) = u_0 + \lambda(u - u_0)$ ;  $u_0$  is a fixed element of  $D(N)$ .

*Proof.* If  $D_t^* = -D_t$  on  $D(N'_u)$  and conditions (2.1) - (2.6) hold, then by Theorem 2.1 it follows that operator (1.1) is potential with respect to (1.7). This implies that we can construct the corresponding functional.

$$F_N[u] = \int_{t_0}^{t_1} \int_0^1 \left[ \left\langle P_{2\tilde{u}(\lambda)} \frac{\partial^2 \tilde{u}(\lambda)}{\partial t^2}, u - u_0 \right\rangle + \left\langle P_{1\tilde{u}(\lambda)} \frac{\partial \tilde{u}(\lambda)}{\partial t}, u - u_0 \right\rangle + \left\langle Q(\tilde{u}(\lambda)), u - u_0 \right\rangle \right] d\lambda dt + F_N[u_0].$$

Let us consider the integral

$$J_2[u] = \int_{t_0}^{t_1} \int_0^1 \left\langle P_{2\tilde{u}(\lambda)} \frac{\partial^2 \tilde{u}(\lambda)}{\partial t^2}, u - u_0 \right\rangle d\lambda dt.$$

Since  $D_t$  is skew-symmetric, we obtain

$$\begin{aligned} J_2[u] &= \int_{t_0}^{t_1} \int_0^1 \left\langle P_{2\tilde{u}(\lambda)} u_{0tt} + \lambda P_{2\tilde{u}(\lambda)} (u_{tt} - u_{0tt}), u - u_0 \right\rangle d\lambda dt \\ &= \int_{t_0}^{t_1} \int_0^1 \left[ \left\langle P_{2\tilde{u}(\lambda)} u_{0tt}, u - u_0 \right\rangle \right. \\ &\quad \left. - \left\langle \lambda (u_t - u_{0t}), \frac{\partial P_{2\tilde{u}(\lambda)}}{\partial t} (u - u_0) + P_{2\tilde{u}(\lambda)} (u_t - u_{0t}) \right\rangle \right] d\lambda dt \\ &= \int_{t_0}^{t_1} \int_0^1 \left[ \left\langle P_{2\tilde{u}(\lambda)} u_{0tt}, u - u_0 \right\rangle + \left\langle \lambda (u - u_0), \frac{\partial^2 P_{2\tilde{u}(\lambda)}}{\partial t^2} (u - u_0) \right\rangle \right. \\ &\quad \left. + \left\langle 2\lambda (u - u_0), \frac{\partial P_{2\tilde{u}(\lambda)}}{\partial t} (u_t - u_{0t}) \right\rangle + \left\langle \lambda (u - u_0), P_{2\tilde{u}(\lambda)} (u_{tt} - u_{0tt}) \right\rangle \right] d\lambda dt \\ &= \int_{t_0}^{t_1} \int_0^1 \left[ \left\langle P_{2\tilde{u}(\lambda)} \frac{\partial^2 \tilde{u}(\lambda)}{\partial t^2}, u - u_0 \right\rangle \right. \\ &\quad \left. + \lambda \left\langle \frac{\partial^2 P_{2\tilde{u}(\lambda)}}{\partial t^2} (u - u_0) + 2 \frac{\partial P_{2\tilde{u}(\lambda)}}{\partial t} (u_t - u_{0t}), u - u_0 \right\rangle \right] d\lambda dt \\ &= \int_{t_0}^{t_1} \int_0^1 \left[ - \left\langle \frac{\partial P_2}{\partial t} (u - u_0), \frac{\partial \tilde{u}(\lambda)}{\partial t} \right\rangle - \left\langle P_2 (u_t - u_{0t}), \frac{\partial \tilde{u}(\lambda)}{\partial t} \right\rangle \right. \\ &\quad \left. + \lambda \left\langle \frac{\partial^2 P_2}{\partial t^2} (u - u_0) + 2 \frac{\partial P_2}{\partial t} (u_t - u_{0t}), u - u_0 \right\rangle \right] d\lambda dt. \end{aligned}$$

For the integral  $J_1[u] = \int_{t_0}^{t_1} \int_0^1 \left\langle P_{1\tilde{u}(\lambda)} \frac{\partial \tilde{u}(\lambda)}{\partial t}, u - u_0 \right\rangle d\lambda dt$  we have

$$\begin{aligned} J_1[u] &= \int_{t_0}^{t_1} \int_0^1 \left[ \left\langle P_{1\tilde{u}(\lambda)} u_{0t}, u - u_0 \right\rangle + \left\langle \lambda P_{1\tilde{u}(\lambda)} (u_t - u_{0t}), u - u_0 \right\rangle \right] d\lambda dt \\ &= \int_{t_0}^{t_1} \int_0^1 \left[ \left\langle P_{1\tilde{u}(\lambda)} u_{0t}, u - u_0 \right\rangle - \left\langle \lambda (u - u_0), D_t (P_{1\tilde{u}(\lambda)}^* (u - u_0)) \right\rangle \right] d\lambda dt. \end{aligned}$$

Further note that

$$D_t[P_{1\tilde{u}(\lambda)}^*(u - u_0)] = D_t \left[ \left( 2 \frac{\partial P_2}{\partial t} - P_{1\tilde{u}(\lambda)} \right) (u - u_0) \right] = 2 \frac{\partial^2 P_2}{\partial t^2} (u - u_0) \\ + 2 \frac{\partial P_2}{\partial t} (u_t - u_{0t}) - \frac{\partial P_{1\tilde{u}(\lambda)}}{\partial t} (u - u_0) - P_{1\tilde{u}(\lambda)} (u_t - u_{0t}) - P'_{1\tilde{u}(\lambda)} \left( u - u_0; \frac{\partial \tilde{u}(\lambda)}{\partial t} \right).$$

Consequently,

$$J_1[u] = \int_{t_0}^{t_1} \int_0^1 \left[ \langle P_{1\tilde{u}(\lambda)} u_{0t}, u - u_0 \rangle - \langle \lambda(u - u_0), 2 \frac{\partial^2 P_2}{\partial t^2} (u - u_0) \rangle \right. \\ \left. - \langle \lambda(u - u_0), 2 \frac{\partial P_2}{\partial t} (u_t - u_{0t}) \rangle + \lambda \langle \frac{\partial P_{1\tilde{u}(\lambda)}}{\partial t} (u - u_0), u - u_0 \rangle \right. \\ \left. + \langle \lambda P_{1\tilde{u}(\lambda)} (u_t - u_{0t}) + \lambda P'_{1\tilde{u}(\lambda)} \left( u - u_0; \frac{\partial \tilde{u}(\lambda)}{\partial t} \right), u - u_0 \rangle \right] d\lambda dt.$$

Taking into consideration (2.12), we get

$$\int_{t_0}^{t_1} \int_0^1 \lambda \langle P'_{1\tilde{u}(\lambda)} \left( u - u_0; \frac{\partial \tilde{u}(\lambda)}{\partial t} \right), u - u_0 \rangle d\lambda dt \\ = \int_{t_0}^{t_1} \int_0^1 \lambda \langle P'_{1\tilde{u}(\lambda)} \left( \frac{\partial \tilde{u}(\lambda)}{\partial t}; u - u_0 \right) - [P'_{1\tilde{u}(\lambda)} \left( \frac{\partial \tilde{u}(\lambda)}{\partial t}; \cdot \right)]^*(u - u_0), u - u_0 \rangle d\lambda dt \\ = \int_{t_0}^{t_1} \int_0^1 \lambda \langle P'_{1\tilde{u}(\lambda)} \left( \frac{\partial \tilde{u}(\lambda)}{\partial t}; u - u_0 \right) - P'_{1\tilde{u}(\lambda)} \left( \frac{\partial \tilde{u}(\lambda)}{\partial t}; u - u_0 \right), u - u_0 \rangle d\lambda dt = 0,$$

i.e.

$$J_1[u] = \int_{t_0}^{t_1} \int_0^1 \left\{ \langle P_{1\tilde{u}(\lambda)} \frac{\partial \tilde{u}(\lambda)}{\partial t}, u - u_0 \rangle + \lambda \langle \frac{\partial P_{1\tilde{u}(\lambda)}}{\partial t} (u - u_0), u - u_0 \rangle \right. \\ \left. - \lambda \langle 2 \frac{\partial^2 P_2}{\partial t^2} (u - u_0), u - u_0 \rangle - \lambda \langle 2 \frac{\partial P_2}{\partial t} (u_t - u_{0t}), u - u_0 \rangle \right\} d\lambda dt.$$

Hence,

$$F_N[u] = \int_{t_0}^{t_1} \int_0^1 \left[ \langle -P_2(u_t - u_{0t}), \frac{\partial \tilde{u}(\lambda)}{\partial t} \rangle - \langle P_{1\tilde{u}(\lambda)}(u - u_0), \frac{\partial \tilde{u}(\lambda)}{\partial t} \rangle \right. \\ \left. + \langle \frac{\partial P_2}{\partial t} (u - u_0), \frac{\partial \tilde{u}(\lambda)}{\partial t} \rangle - \lambda \langle \frac{\partial^2 P_2}{\partial t^2} (u - u_0), u - u_0 \rangle \right]$$

$$+ \left\langle Q(\tilde{u}(\lambda)) + \lambda \frac{\partial P_{1\tilde{u}(\lambda)}}{\partial t}(u - u_0), u - u_0 \right\rangle d\lambda dt + F_N[u_0].$$

Using (2.15) - (2.17), we get

$$F_N[u] = \int_{t_0}^{t_1} \left[ \langle \mathcal{R}_2 u_t, u_t \rangle + \langle \mathcal{R}_1(u) - \frac{\partial \mathcal{R}_2}{\partial t} u, u_t \rangle + \mathcal{B}[u] \right] dt + F_N[u_0]. \quad (2.18)$$

It is easy to show that

$$\begin{aligned} \delta F_N[u, h] &= \int_{t_0}^{t_1} \left[ \langle \mathcal{R}_2 h_t + \mathcal{R}'_{1u} h, u_t \rangle + \langle \mathcal{R}_2 u_t + \mathcal{R}_1(u), h_t \rangle \right. \\ &\quad \left. - \langle \frac{\partial \mathcal{R}_2}{\partial t} h, u_t \rangle - \langle \frac{\partial \mathcal{R}_2}{\partial t} u, h_t \rangle + \langle \text{grad}_{\Phi_1} \mathcal{B}[u], h \rangle \right] dt \\ &= \int_{t_0}^{t_1} \left[ \langle -\frac{\partial \mathcal{R}_2^*}{\partial t} u_t, h \rangle - \langle \mathcal{R}_2^* u_{tt}, h \rangle - \langle \frac{\partial \mathcal{R}_2}{\partial t} u_t, h \rangle - \langle \mathcal{R}_2 u_{tt}, h \rangle \right. \\ &\quad \left. + \langle \mathcal{R}'_{1u} u_t, h \rangle - \langle \frac{\partial \mathcal{R}_1}{\partial t}(u), h \rangle - \langle \mathcal{R}'_{1u} u_t, h \rangle - \left( \frac{\partial \mathcal{R}_2}{\partial t} \right)^* u_t, h \rangle \right. \\ &\quad \left. + \langle \frac{\partial^2 \mathcal{R}_2}{\partial t^2} u, h \rangle + \langle \frac{\partial \mathcal{R}_2}{\partial t} u_t, h \rangle + \langle \text{grad}_{\Phi_1} \mathcal{B}[u], h \rangle \right] dt \\ &= \int_{t_0}^{t_1} \langle (-\mathcal{R}_2 - \mathcal{R}_2^*) u_{tt} + (\mathcal{R}'_{1u} - \mathcal{R}'_{1u} - 2 \frac{\partial \mathcal{R}_2^*}{\partial t}) u_t \\ &\quad + \left( \text{grad}_{\Phi_1} \mathcal{B}[u] - \frac{\partial \mathcal{R}_1}{\partial t}(u) + \frac{\partial^2 \mathcal{R}_2}{\partial t^2} u \right), h \rangle dt = \int_{t_0}^{t_1} \langle N(u), h \rangle dt \end{aligned}$$

$$\forall u \in D(N), \forall h \in D(N'_u).$$

Thus, the sufficiency of representation (2.14) is proved. On the other hand, if  $D_t^* = -D_t$  on  $D(N'_u)$  and equation (1.1) can be represented in the form (2.14), then

$$P_2 = -\mathcal{R}_2 - \mathcal{R}_2^*, \quad (2.19)$$

$$P_{1u} = \mathcal{R}'_{1u} - \mathcal{R}'_{1u} - 2 \frac{\partial \mathcal{R}_2^*}{\partial t}, \quad (2.20)$$

$$Q(u) = \text{grad}_{\Phi_1} \mathcal{B}[u] - \frac{\partial \mathcal{R}_1}{\partial t}(u) + \frac{\partial^2 \mathcal{R}_2}{\partial t^2} u. \quad (2.21)$$

□

**Remark 2.** If  $0_U \in D(N)$ , then the operators  $\mathcal{R}_1$ ,  $\mathcal{R}_2$  can be found by formulas

$$\langle \mathcal{R}_1(u), u_t \rangle = \int_0^1 \langle -\lambda P_{1\lambda} u, u_t \rangle d\lambda,$$

$$\langle \mathcal{R}_2 u_t, u_t \rangle = \int_0^1 \langle -\lambda P_2 u_t, u_t \rangle d\lambda.$$

In the case of bilinear mapping (1.10) we have

$$\mathcal{R}_1(u) = \int_0^1 -\lambda P_{1\lambda} u d\lambda,$$

$$\mathcal{R}_2 = -\frac{1}{2} P_2,$$

**Example.** Consider the following partial differential equation:

$$N(u) \equiv u_{tt} + 2\beta v(t)u_{tx} + u_{xxxx} + v^2(t)u_{xx} + \beta v'(t)u_x = 0, \quad (2.22)$$

$$(x, t) \in Q_T = (a, b) \times (0, T),$$

where  $\beta$  is a constant,  $v(t)$  is a fixed function and  $u(x, t)$  is the unknown function.

We denote by  $D(N)$  the domain of definition of the operator  $N$  in (2.22):

$$D(N) = \{u \in U = C_{t,x}^{2,4}(\overline{Q_T}) : u|_{t=0} = \phi_1(x), u_t|_{t=0} = \phi_2(x) \ (x \in (a, b)), \quad (2.23)$$

$$u|_{x=a} = \psi_1(t), \ u|_{x=b} = \psi_2(t) \ (t \in (0, T))\},$$

where  $\phi_i, \psi_i$  ( $i=1,2$ ) are given functions.

We denote  $V = C(\overline{Q_T})$  and determine the bilinear form  $\Phi(\cdot, \cdot) : V \times U \rightarrow \mathbb{R}$  by setting

$$\Phi(v, g) = \int_0^T \int_a^b v(x, t)g(x, t) dx dt. \quad (2.24)$$

Let us note that equation (2.22) has the structure of equation (1.1). Indeed, in this case

$$P_2 = I, \quad P_1 = 2\beta v(t)D_x, \quad Q(u) = u_{xxxx} + v^2(t)u_{xx} + \beta v'(t)u_x.$$

Let us show that conditions (2.1) – (2.6) are satisfied. We have

$$(2.1) \implies I - I = 0,$$

$$(2.2) \implies 0 = 0,$$

$$(2.3) \implies 2\beta v(t)D_x - 2\beta v(t)D_x = 0,$$

$$(2.4) \implies -2\beta v'(t)D_x + D_x^4 + v^2(t)D_x^2 + \beta v'(t)D_x - D_x^4 - v^2(t)D_x^2 + \beta v'(t)D_x = 0,$$

$$(2.5) \implies 0 = 0,$$

$$(2.6) \implies 0 = 0.$$

Hence, equation (2.22) can be represented in the form of the Euler-Lagrange equation. Further, using (2.15) – (2.17) one obtains

$$\mathcal{R}_2 = -\frac{1}{2}I, \quad \mathcal{R}_1 = -\beta v(t)D_x, \quad \mathcal{B}[u] = \frac{1}{2} \int_a^b \{u_{xx}^2 - v^2(t)u_x^2\} dx.$$

Thus, the potential of the operator  $N$  in (2.22) can be written in the form

$$F_N[u] = \frac{1}{2} \int_0^T \int_a^b \{-u_t^2 - 2\beta v(t)u_x u_t + u_{xx}^2 - v^2(t)u_x^2\} dx dt + F_N[u_0]. \quad (2.25)$$

Let us note that potential (2.25) can be used to obtain an infinite number of first integrals of the given equation (2.22).

Suppose that  $u \in C_{t,x}^{2,\infty}(\overline{Q_T})$ , then the first integrals of (2.22) are

$$I_k[u] = \int_a^b (-u_t - \beta v(t)u_x) \frac{\partial^{2k+1}u}{\partial x^{2k+1}} dx, \quad (k = 0, 1, 2, \dots)$$

(see [7]).

This paper can be considered as a continuation of [4].

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